# Department of Mathematics, University of Utah <br> Applied Complex Variables and Asymptotic Methods <br> MATH 6720 - Section 001 - Spring 2024 <br> Homework 4 Solutions <br> Complex Integration, II 

Due: Friday, Feb 23, 2024

Below, problem C in section A.B is referred to as exercise A.B.C.
Text: Complex Variables: Introduction and Applications, Ablowitz \& Fokas,
Exercises: 3.3.3
3.3.4
3.5.1, parts a) - f)
3.5.2, parts a) - c)

Submit your homework assignment on Canvas via Gradescope.
3.3.3. Given the function

$$
f(z)=\frac{z}{(z-2)(z+i)},
$$

expand $f(z)$ in a Laurent series in powers of $z$ in the regions,
(a) $|z|<1$
(b) $1<|z|<2$
(c) $|z|>2$

Solution: Before delving into the individual parts of this problem, we perform some preliminary computations. First, expanding $f$ in partial fractions yields the representation,

$$
\begin{equation*}
f(z)=\frac{\frac{2}{2+i}}{z-2}+\frac{\frac{i}{2+i}}{z+i}, \tag{1}
\end{equation*}
$$

so that we can accomplish this problem by considering Laurent series' for $1 /(z-2)$ and $1 /(z+i)$. For the first expansion, we note:

$$
\begin{aligned}
& |z|<2: \frac{1}{z-2}=\left(-\frac{1}{2}\right) \frac{1}{1-(z / 2)} \stackrel{|z / 2|<1}{=}-\frac{1}{2} \sum_{j=0}^{\infty}\left(\frac{z}{2}\right)^{j}=\sum_{j=0}^{\infty}-\frac{1}{2^{j+1}} z^{j}, \\
& |z|>2: \frac{1}{z-2}=\left(\frac{1}{z}\right) \frac{1}{1-(2 / z)} \stackrel{|2 / z|<1}{=} \frac{1}{z} \sum_{j=0}^{\infty}\left(\frac{2}{z}\right)^{j}=\sum_{j=-\infty}^{-1} \frac{1}{2^{j+1}} z^{j},
\end{aligned}
$$

and this implies that,

$$
\frac{1}{z-2}=\left\{\begin{array}{ll}
\sum_{j=0}^{\infty} a_{j} z^{j}, & |z|<2,  \tag{2}\\
\sum_{j=-\infty}^{-1} a_{j} z^{j}, & |z|>2,
\end{array} \quad a_{j}= \begin{cases}-\frac{1}{2^{j+1}}, & j \geq 0 \\
\frac{1}{2^{j+1}}, & j<0\end{cases}\right.
$$

Similarly, for $1 /(z+i)$ we have,

$$
\begin{aligned}
& |z|<1: \frac{1}{z+i}=\left(\frac{1}{i}\right) \frac{1}{1-i z} \stackrel{|i z|<1}{=} \frac{1}{i} \sum_{j=0}^{\infty}(i z)^{j}=\sum_{j=0}^{\infty} i^{j-1} z^{j}, \\
& |z|>1: \frac{1}{z+i}=\left(\frac{1}{z}\right) \frac{1}{1-(1 / i z)} \stackrel{|i z|>1}{=} \frac{1}{z} \sum_{j=0}^{\infty} \frac{1}{(i z)^{j}}=\sum_{j=-\infty}^{-1} i^{j+1} z^{j},
\end{aligned}
$$

and so,

$$
\frac{1}{z+i}=\left\{\begin{array}{ll}
\sum_{j=0}^{\infty} b_{j} z^{j}, & |z|<1,  \tag{3}\\
\sum_{j=-\infty}^{-1} b_{j} z^{j}, & |z|>1,
\end{array} \quad b_{j}= \begin{cases}i^{j-1}, & j \geq 0 \\
i^{j+1}, & j<0\end{cases}\right.
$$

We can now answer the individual parts of this question:
(a) For $|z|<1$, we use the Taylor series both from (2) and (3) in (1) to obtain,

$$
f(z)=\sum_{j=0}^{\infty} c_{j} z^{j}, \quad c_{j}=\frac{2 a_{j}}{2+i}+\frac{i b_{j}}{2+i}=\frac{1}{2+i}\left(-\frac{1}{2^{j}}+i^{j}\right), j \geq 0 .
$$

(b) For $1<|z|<2$, we use the Laurent series from (2) and the Taylor series from (3), obtaining,

$$
f(z)=\sum_{j=-\infty}^{\infty} c_{j} z^{j}, \quad c_{j}=\left\{\begin{array}{cc}
\frac{2 a_{j}}{2+i}=-\frac{1}{(2+i) 2^{j}}, & j \geq 0, \\
\frac{i b_{j}}{2+i}=\frac{i+2+2}{2+i}, & j<0
\end{array}\right.
$$

(c) For $|z|>2$, we exercise the Laurent series from both (2) and (3) to obtain,

$$
f(z)=\sum_{j=-\infty}^{-1} c_{j} z^{j}, \quad c_{j}=\frac{2 a_{j}}{2+i}+\frac{i b_{j}}{2+i}=\frac{1}{2+i}\left(\frac{1}{2^{j}}-i^{j}\right), j<0
$$

3.3.4. Evaluate the integral $\oint_{C} f(z) \mathrm{d} z$ where $C$ is the unit circle centered at the origin and $f(z)$ is given as follows,
(a) $\frac{e^{z}}{z^{3}}$
(b) $\frac{1}{z^{2} \sin z}$
(c) $\tanh z$
(d) $\frac{1}{\cos ^{2 z}}$
(e) $e^{1 / z}$

Solution: Assuming $f$ is analytic in some open domain of $|z|=1$, our strategy in this problem is to utilize the Laurent series formula,

$$
g(z)=\sum_{n \in \mathbb{Z}} c_{n} z^{n}, \quad c_{n}=\frac{1}{2 \pi i} \oint_{C} \frac{g(z)}{z^{n+1}} \mathrm{~d} z
$$

so that, by appropriately defining $g$, integrals around $C$ can be computed by identifying certain coefficients of the Laurent series of $g$. (Note that this is equivalent to expanding $f$ in a Laurent series and computing the $c_{-1}$ coefficient for $f$.)
(a) With $g(z)=e^{z}$, then

$$
\oint_{C} f(z) \mathrm{d} z=\oint_{C} \frac{g(z)}{z^{3}} \mathrm{~d} z=2 \pi i c_{2} .
$$

The function $g$ has the convergent Taylor series $g(z)=1+z+z^{2} / 2+\cdots$, so we immediately conclude that $c_{2}=1 / 2$, and thus, $\oint_{C} f(z) \mathrm{d} z=\pi i$.
(b) We take $g(z)=z /(\sin z)$, so that,

$$
\oint_{C} f(z) \mathrm{d} z=\oint_{C} \frac{g(z)}{z^{3}} \mathrm{~d} z=2 \pi i c_{2} .
$$

The Laurent series expansion around $z=0$ for $g$, which is analytic around $|z|=1$, can be computed via,

$$
\begin{aligned}
\frac{z}{\sin z}=\frac{z}{z-\frac{z^{3}}{3!}+\frac{z^{5}}{5!}-\cdots}=\frac{1}{1-\frac{z^{2}}{3!}+\frac{z^{4}}{5!}-\cdots} & =\frac{1}{1-\left(\frac{z^{2}}{3!}-\frac{z^{4}}{5!}+\cdots\right)} \\
& =1+\left(\frac{z^{2}}{3!}-\frac{z^{4}}{5!}+\cdots\right)+\cdots,
\end{aligned}
$$

where the last equality uses the geometric series on the terms in parenthesis. From this expression, we see that $c_{2}=1 / 3!=1 / 6$, so that $\oint_{C} f(z) \mathrm{d} z=\pi i / 3$.
(c) The function $\tanh z$ is analytic everywhere except where the denominator $(\cosh z)$ vanishes. The denominator vanishes at $z=(2 j+1) i \pi / 2$, for $j \in \mathbb{Z}$. Note that $C$ does not enclose any of these points, and so $f$ is analytic inside $C$. By the Cauchy-Goursat theorem, $\oint_{C} f(z) \mathrm{d} z=0$.
(d) For this problem, we cannot directly compute a Laurent series since $\cos 2 z=0$ for $z= \pm \frac{\pi}{4}$, which both lie within $C$. Instead, we deform $C$ into two circles of small radius (say radius smaller than $\frac{\pi}{4}$ ) with a crosscut connecting them, with the first circle $C_{-}$ counterclockwise around $z=-\pi / 4$, and second circle $C_{+}$counterclockwise around $\pi / 4$. The two crosscut integrals add-out, and so we must compute,

$$
\int_{C} \frac{1}{\cos 2 z} \mathrm{~d} z=\int_{C_{-}} \frac{1}{\cos 2 z} \mathrm{~d} z+\int_{C_{+}} \frac{1}{\cos 2 z} \mathrm{~d} z .
$$

For the $C_{-}$integral, $g_{-}(z):=\frac{1}{\cos 2 z}$ is analytic in a neighborhood of $z=-\pi / 4$, and we have,
$g_{-}(z)=\frac{1}{\cos 2 z}=\frac{1}{\sin (2(z+\pi / 4))}=\frac{1}{2(z+\pi / 4)} \sum_{j \geq 0}\left[\frac{(2(z+\pi / 4))^{2}}{3!}-\frac{(2(z+\pi / 4))^{4}}{5!}-\ldots\right]^{j}$
where the last equality is the Laurent series of $1 / \sin w$ around $w=0$ with $w \leftarrow 2(z+\pi / 4)$. (The computation uses the Taylor expansion of $\sin w$ along with a geometric series.) Hence, the Laurent series coefficient $c_{-1}$ for $g_{-}$around $z=-\pi / 4$ is given by $c_{-1}=\frac{1}{2}$, so that,

$$
\int_{C_{-}} \frac{1}{\cos 2 z} \mathrm{~d} z=\frac{1}{2} .
$$

A similar computation with,
$g_{+}(z):=\frac{1}{\cos 2 z}=-\frac{1}{\sin (2(z-\pi / 4))}=\frac{-1}{2(z-\pi / 4)} \sum_{j \geq 0}\left[\frac{(2(z-\pi / 4))^{2}}{3!}-\frac{(2(z-\pi / 4))^{4}}{5!}-\ldots\right]^{j}$,
implies that

$$
\int_{C_{+}} \frac{1}{\cos 2 z} \mathrm{~d} z=-\frac{1}{2} .
$$

Hence,

$$
\int_{C} \frac{1}{\cos 2 z} \mathrm{~d} z=\int_{C_{-}} \frac{1}{\cos 2 z} \mathrm{~d} z+\int_{C_{+}} \frac{1}{\cos 2 z} \mathrm{~d} z=\frac{1}{2}-\frac{1}{2}=0
$$

(e) Again defining $g=f$, we have

$$
\oint_{C} f(z) \mathrm{d} z=2 \pi i c_{-1} .
$$

The Laurent series for $f=e^{1 / z}$ is given by,

$$
e^{1 / z} \stackrel{z=1 / w}{=} e^{w}=\sum_{j \geq 0} \frac{w^{j}}{j!} \stackrel{z=1 / w}{=} \sum_{j \leq 0} \frac{z^{j}}{(-j)!}=1+\frac{1}{z}+\frac{1}{2 z^{2}}+\cdots,
$$

so that $c_{-1}=1$, and hence $\oint_{C} f(z) \mathrm{d} z=2 \pi i$.
3.5.1. Discuss the type of singularity (removable, pole and order, essential, branch, cluster, natural barrier, etc.); if the type is a pole give the strength of the pole, and give the nature (isolated or not) of all singular points associated with the following functions. Include the point at infinity.
(a) $\frac{e^{z^{2}}-1}{z^{2}}$
(b) $\frac{e^{2 z}-1}{z^{2}}$
(c) $e^{\tan z}$
(d) $\frac{z^{3}}{z^{2}+z+1}$
(e) $\frac{z^{1 / 3}-1}{z-1}$
(f) $\log \left(1+z^{1 / 2}\right)$
(i) $\operatorname{sech} z$
(j) $\operatorname{coth} 1 / z$

## Solution:

(a) The point $z=0$ is a removable (hence isolated) singularity, and $z=\infty$ is an essential (hence isolated) singularity. The finite- $z$ singularities only occur when the denominator
vanishes, i.e., $z=0$, where $f$ (and also $f^{\prime}$ ) are undefined. Through a Taylor expansion of the numerator, we find,

$$
\frac{e^{z^{2}}-1}{z^{2}}=\frac{1}{z^{2}}\left[-1+\sum_{j=0}^{\infty} \frac{z^{2 j}}{j!}\right]=\sum_{j=0}^{\infty} \frac{z^{2 j}}{(j+1)!},
$$

hence this function has a convergent power series around $z=0$, and thus $z=0$ is a removable singularity. Redefining the function to have value 1 at $z=0$ makes it analytic. To investigate $z=\infty$, consider the expansion above in the variable $w=1 / z$. Then as a function of $w$ this function has an infinite principal part of its Laurent expansion around $w=0$, so that $w=0$ (hence $z=\infty$ ) is an essential singularity.
(b) The point $z=0$ is a simple pole (hence isolated) with strength 2 , and $z=\infty$ is an essential (isolated) singularity. Again, we take a Taylor series expansion and simplify:

$$
\frac{e^{2 z}-1}{z^{2}}=\frac{1}{z^{2}}\left[-1+\sum_{j=0}^{\infty} \frac{(2 z)^{j}}{j!}\right]=\frac{2}{z}+2+\ldots,
$$

where $\ldots$ is a power series around $z=0$. Thus, $z=0$ is a simple pole of strength 2. Making the definition $w=1 / z$, we see that the again the the principal part of the expansion above in terms of $w$ has an infinite number of terms, so that $w=0$ (hence $z=\infty)$ is an essential singularity.
(c) The point $z=\infty$ is a cluster (non-isolated) singular point, and $z=\frac{(2 n+1) \pi}{2}$ for $n \in \mathbb{Z}$ are essential (isolated) singularities. Making the substituion $w=1 / z$, then $\exp (\tan 1 / w)$ has singularities for $w=\frac{2}{(2 n+1)}, n \in \mathbb{Z}$, which cluster around 0 , so $w=0(z=\infty)$ is a cluster singularity. To establish that the finite $z$ points are essential singularities, consider $z=-\pi / 2$, and note that in a neighborhood around this point,
$\exp (\tan z)=\exp \left(\sin z \frac{1}{\sin (z+\pi / 2)}\right)=\exp \left(\frac{\sin z}{z+\pi / 2} \sum_{j \geq 0}\left[\frac{(z+\pi / 2)^{2}}{3!}-\frac{(z+\pi / 2)^{4}}{5!}+\ldots\right]^{j}\right)$,
and hence around $z=-\pi / 2$, this function behaves like $\exp (1 / z)$ around $z=0$, i.e., as an essential singularity.
(d) The point $z=-\frac{1}{2}+i \frac{\sqrt{3}}{2}$ is a simple pole (isolated) with strength $-i / \sqrt{3}$, the point $z=\frac{1}{2}+i \frac{\sqrt{3}}{2}$ is a simple pole (isolated) with strength $i / \sqrt{3}$, and $z=\infty$ is a simple pole (isolated) of strength 1 . The point at $z=\infty$ is the simplest to see with $w=1 / z$ and since for large $z, \frac{z^{3}}{z^{2}+z+1} \sim z=\frac{1}{w}$, which is a simple pole of strength 1 . The other two poles are found at the roots of the denominator:

$$
z^{2}+z+1=0 \Longrightarrow z=z_{0}, \bar{z}_{0}=-\frac{1}{2}+i \frac{\sqrt{3}}{2},-\frac{1}{2}-i \frac{\sqrt{3}}{2}
$$

Hence, the function can be rewritten as,

$$
\frac{z^{3}}{z^{2}+z+1}=\frac{z^{3}}{\left(z-z_{0}\right)\left(z-\bar{z}_{0}\right)}
$$

corresponding to simple poles at $z=z_{0}, \bar{z}_{0}$. The strength of the pole at $z=z_{0}$ is given by,

$$
\frac{z_{0}^{3}}{z_{0}-\bar{z}_{0}}=-\frac{i}{\sqrt{3}} z_{0}^{3}=-\frac{i}{\sqrt{3}} z_{0}^{2} z_{0} \stackrel{z_{0}^{2}+z_{0}+1=0}{=}-\frac{i}{\sqrt{3}}\left(-1-z_{0}\right) z_{0}=-\frac{i}{\sqrt{3}}\left(-z_{0}-z_{0}^{2}\right)=-\frac{i}{\sqrt{3}}
$$

A similar computation can be carried out for the pole at $\bar{z}_{0}$, showing its strength is $+i / \sqrt{3}$.
(e) The points $z=0, \infty$ are each a branch (non-isolated) singularity. For finite $z$, this function has three branches. For two of the branches, $z=1$ is a simple pole (isolated) with strengths $-\frac{3}{2}+i \sqrt{3} / 2$ and $-\frac{3}{2}-i \sqrt{3} / 2$, respectively. The third branch has a removable (isolated) singularity at $z=1$. To see the branch point at $\infty$, make the substitution $w=1 / z$, so that,

$$
\frac{z^{1 / 3}-1}{z-1}=\frac{w\left(w^{-1 / 3}-1\right)}{1-w} .
$$

Thus, we see through the $w^{-1 / 3}$ term that $w=0(z=\infty)$ is a branch point singularity for this function. The point $z=0$ is a branch point as can be seen directly by definition of the function. To assess the situation around $z=1$, we make the branch cut of the function along the negative real axis, corresponding to three branches:

- Branch 1: $z^{1 / 3}=|z| e^{i(\arg z) / 3}$ for $\arg z \in[-\pi, \pi)$
- Branch 2: $z^{1 / 3}=|z| e^{i(\arg z) / 3}$ for $\arg z \in[\pi, 3 \pi)$
- Branch 3: $z^{1 / 3}=|z| e^{i(\arg z) / 3}$ for $\arg z \in[3 \pi, 5 \pi)$

On branch 1 around $z=1$, we have $z^{1 / 3} \rightarrow z^{1 / 3}$ for $\arg z$ around 0 . Using the Taylor expansion of $z^{1 / 3}$ around $z=1$, we have,

$$
\frac{-1+z^{1 / 3}}{z-1}=\frac{-1+1+\frac{1}{3}(z-1)+\ldots}{z-1}=\frac{1}{3}+\ldots,
$$

where $\ldots$ denote higher-order power series terms. Hence, $z=1$ for Branch 1 a removable singularity, and defining the function to have value $1 / 3$ at $z=1$ makes it analytic around $z=1$. On branch 2, we have $z^{1 / 3} \rightarrow e^{i 2 \pi / 3} z^{1 / 3}=\left(-\frac{1}{2}+i \frac{\sqrt{3}}{2}\right) z^{1 / 3}$ for $\arg z$ around 0 . In this case, the numerator is analytic around $z=1$ but does not vanish there (instead taking on value $-3 / 2+i \sqrt{\frac{3}{2}}$ ), so we conclude that $z=1$ is a simple pole with strength $-3 / 2+i \sqrt{\frac{3}{2}}$. A similar argument can be made for branch 3 .
(f) The points $z=0, \infty$ are branch (non-isolated) points, and $z=1$ is a branch (nonisolated) point as well. That $z=0, \infty$ are branch points is a direct result from the fact that these are branch points for $z \mapsto \sqrt{z}$. To see that $z=1$ is a branch point, consider the (non-principal) branch of $\sqrt{z}$, given by $-\sqrt{z}$ for $\arg z \in[-\pi, \pi)$. Then,

$$
\log (1+\sqrt{z}) \rightarrow \log (1-\sqrt{z})
$$

so that $z=1$ is also a branch point since $\log w$ has a branch point at $w=0$. No additional branch points are given by the principal branch of $\sqrt{z}$ since $1+\sqrt{z}$ cannot vanish on the principal branch of $\sqrt{z}$.
3.5.2. Evaluate the integral $\oint_{C} f(z) \mathrm{d} z$, where $C$ is a unit circle centered at the origin, and where $f(z)$ is given below.
(a) $\frac{g(z)}{z-w}, g(z)$ entire
(b) $\frac{z}{z^{2}-w^{2}}$
(c) $z e^{1 / z^{2}}$

Solution: We will use characterizations of singularities to address these problems, although an equally acceptable (and frequently simpler) way is to use the Cauchy integral formula. We recall from the solution of problem 3.3.4 that if $f$ has a Laurent series around 0 , then the integral of $f$ around a simple closed contour $C$ lying within the region of convergence of the series equals the $c_{-1}$ coefficient of the series. In addition, we assume in what follows that $w$ either is enclosed by $C$, or lies outside it. (If $w$ lies on $C$, then the integrals do not have convergent values.)
(a) Since $g$ is entire, then we may expand it in the Taylor series,

$$
g(z)=\sum_{j \geq 0} \frac{g^{(j)}(w)}{j!}(z-w)^{j} .
$$

which is uniformly convergent on and in a neighborhood of $C$. Then using the fact that for an arbitrary $k \in \mathbb{Z}, \oint_{C}(z-w)^{k} \mathrm{~d} z=2 \pi i$ if and only if both $k=-1$ and $C$ encloses $w$ (and is zero otherwise), then

$$
\begin{aligned}
& \oint_{C} \frac{g(z)}{z-w} \mathrm{~d} z=\sum_{j \geq 0} \frac{g^{(j)}(w)}{j!} \oint_{C}(z-w)^{j-1} \mathrm{~d} z=g(w) \oint_{C}(z-w)^{-1} \mathrm{~d} z \\
&=\left\{\begin{array}{r}
2 \pi i g(w), \\
0,
\end{array} \text { otherwise } w\right. \\
& 0,
\end{aligned}
$$

(b) If $w$ (and hence also $-w$ ) lies outside $C$, then $f(z)$ is only singular at points outside $C$, and therefore is analytic inside $C$. Then by Cauchy-Goursat, $\oint_{C} f(z) \mathrm{d} z=0$. If instead $w$ is enclosed by $C$, then we use partial fractions along with a Laurent series expansion around 0 using the geometric series, valid for $|z|>|w|$ and thus in particular on $C$ :

$$
\begin{aligned}
\frac{z}{z^{2}-w^{2}}=\frac{1 / 2}{z-w}+\frac{1 / 2}{z+w} & =\left(\frac{1}{2 z}\right) \frac{1}{1-\frac{w}{z}}+\left(\frac{1}{2 z}\right) \frac{1}{1+\frac{w}{z}} \\
& =\frac{1}{2 z} \sum_{j \geq 0}\left(\frac{w}{z}\right)^{j}+\frac{1}{2 z} \sum_{j \geq 0}\left(-\frac{w}{z}\right)^{j} .
\end{aligned}
$$

Again we have $\oint_{C} z^{k} \mathrm{~d} z=0$ for $k<-1$, so only the $j=0$ terms (corresponding to simple poles of $f$ ) produce non-zero contribution by integrating over $C$ :

$$
\oint_{C} \frac{z}{z^{2}-w^{2}} \mathrm{~d} z=\oint_{C} \frac{1}{2 z} \mathrm{~d} z+\oint_{C} \frac{1}{2 z} \mathrm{~d} z=2 \pi i .
$$

(c) The function $e^{1 / z^{2}}$ has an essential singularity at $z=0$. Using a Laurent expansion around 0 (valid on $C$ ), then:

$$
\oint_{C} z e^{1 / z^{2}} \mathrm{~d} z=\oint_{C} \sum_{j \geq 0} \frac{z^{1-2 j}}{j!} \mathrm{d} z=\sum_{j \geq 0} \frac{1}{j!} \oint_{C} z^{1-2 j} \mathrm{~d} z=\oint_{C} z^{-1} \mathrm{~d} z=2 \pi i
$$

