

## 2 **Enhancing Understand of Particle Simulations Through** 3 **Deformation-Based Visualization**

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6 **Abstract:** We present a physically based method for visualizing deformation in  
7 particle simulations, such as those describing structural mechanics simulations.  
8 The method uses the *deformation gradient* tensor to transform carefully chosen  
9 glyphs representing each particle. The visualization approximates how simulated  
10 objects responding to applied forces might look in reality, allowing for a better  
11 understanding of material deformation, an important indicator of, for example, ma-  
12 terial failure. It can also help highlight possible errors and numerical deficiencies in  
13 the simulation itself, suggesting how simulations might be changed to yield more  
14 accurate results.

15 **Keywords:** Deformation, visualization, particle methods, material point method.

### 16 **1 Introduction and Background**

17 Particle methods [Belytschko, Krongauz, Organ, Fleming, and Krysl (1996)] are  
18 numerical simulation methods in which materials are modeled by collections of  
19 discrete computational particles, which can move about the computational domain  
20 as indicated by the model equations. Generally speaking, these methods produce  
21 output in which each particle is identified by its location and additional data values.  
22 The values can be of any type, and methods producing scalar, vector, and tensor  
23 values are common. Compared to the well-known finite element method (FEM)  
24 [Bathe (1996)], particle methods have the major advantage that they are well-suited  
25 to handling large deformations, such as might be found in simulations leading to  
26 material failure, for example. When dealing with such deformations, FEM can  
27 suffer from entangled, inverted, or otherwise ill-conditioned meshes; remeshing  
28 can alleviate these problems but only at heavy cost, both in terms of performance

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29 and accuracy. For simulations in which large deformations are expected, as in  
30 biomechanics, or where typical engineering materials will experience failure, such  
31 as explosions, particle methods, which are not affected by these concerns, can be a  
32 suitable alternative.

33 Because the particles “carry” data values, a glyph-based approach to visualization  
34 that places geometry at the position of each particle is suitable. Scalar data val-  
35 ues can be displayed by varying free parameters, such as glyph size and color.  
36 However, it is not clear how to visualize deformation within this framework, an im-  
37 portant quantity in structural mechanics that conveys information about the strains  
38 and resulting stresses experienced by an object of study.

39 Deformation is an important physical response to loading, and an independent vari-  
40 able in models of material failure [Maloney and Lemaître (2004); Meakin (1991)];  
41 as such it is an important consideration in simulations involving large deforma-  
42 tions, such as for penetrative tissue damage [Ionescu, Guilkey, Berzins, Kirby, and  
43 Weiss (2006)], high-temperature damage in pipes [Hall and Hayhurst (1991)], and  
44 explosion in containers of high-energy materials [Guilkey, Harman, and Banerjee  
45 (2007)]. Because deformation is directly observable by humans in physical ob-  
46 jects, it should be included in visualization of such simulations, at the very least as  
47 an indirect way to validate simulation data. Leaving deformation out of the process  
48 produces at best a deficient visualization, and at worst, a misleading one.

49 In their common usage, glyph-based visualization approaches force users to infer  
50 deformation from the relative motion of the glyphs, and for many arrangements,  
51 the actual deformation can be difficult to discern. This paper presents a particular  
52 approach to glyph-based visualization for particle data that uses carefully chosen  
53 glyphs and the *deformation gradient* to clearly display deformation, both at local  
54 and global scales. We demonstrate how the method improves on current techniques,  
55 allowing scientists a better understanding of their simulations.

## 56 **1.1 Particle Methods**

57 Particle methods form a subset of the mesh free [Liu (2003)] or meshless [Be-  
58 lytschko, Krongauz, Organ, Fleming, and Krysl (1996); Atluri, Liu, and Han (1998)]  
59 methods, which differ from FEM (and other fully meshed methods such as finite  
60 differences, etc.) in that no object geometry is represented by a mesh. Particle  
61 methods represent objects by discretizing them into collections of particles, each  
62 of which is a Lagrangian representation of some part of the object either directly  
63 as a small continuum of matter, or a sampling point. Particles are, in principle, free  
64 to move independently about the domain, but they are usually restricted to realistic  
65 behavior by material models and other physical constraints.

66 Examples of particle methods include smoothed particle hydrodynamics (SPH)  
 67 [Monaghan (2005)], the related smooth particle applied mechanics [Hoover (2006)],  
 68 and the particle-in-cell (PIC) family of methods [Harlow (1964); Brackbill and  
 69 Ruppel (1986)], in which a stationary, background mesh is used to compute gradi-  
 70 ents and the particles move through the grid cells as the simulation proceeds. The  
 71 material point method (MPM) [Sulsky, Chen, and Shreyer (1994); Sulsky, Zhou,  
 72 and Schreyer (1995)] (along with its successor, the generalized interpolation ma-  
 73 terial point method (GIMP) [Bardenhagen and Kober (2004)]) is a PIC method  
 74 explicitly designed for performing structural mechanics simulations. MPM rep-  
 75 represents geometry by discretizing objects into particles, or *material points*. Each  
 76 material point in such a particle model represents a small piece of material from  
 77 the object and obeys the laws of continuum mechanics. Each material point carries  
 78 several physical parameters, such as mass, volume, stress, etc. In the simulation, the  
 79 particles move as the object responds to loads placed on it, resulting in an overall  
 80 deformation of the object. In the process of updating particle volumes, the MPM al-  
 81 gorithm may compute a tensor quantity known as the *deformation gradient*, which  
 82 acts as a local measure of distortion affecting each particle. This value may be used  
 83 as input to the material constitutive model, which relates stress to strain; however  
 84 our method takes advantage of the deformation gradient to visualize deformation.

## 85 1.2 Overview of MPM/GIMP

86 Numerous flavors exist of both standard MPM and GIMP. Particular choices of  
 87 grid and particle basis functions [Steffen, Wallstedt, Guilkey, Kirby, and Berzins  
 88 (2008)], whether or not to lump the mass matrix [Love and Sulsky (2006)], and  
 89 what time-stepping algorithm to use (implicit, explicit, etc.) [Guilkey and Weiss  
 90 (2003); Wallstedt and Guilkey (2008)] will have large impacts on implementation  
 91 details as well as expected algorithm performance. Here, we present a very brief  
 92 overview of a single timestep in the Update Stress Last (USL) explicit method. We  
 93 do this to illustrate the role that the deformation gradient plays in the method. For  
 94 full implementation details, we suggest referring to Wallstedt and Guilkey (2008).

The simulation begins with a collection of particles approximating the geometry  
 of the objects of interest (Figure 1). Each particle is assigned various initial quan-  
 tities: mass ( $m_p$ ), position ( $\mathbf{x}_p^0$ ), velocity ( $\mathbf{v}_p^0$ ), deformation gradient ( $\mathbf{F}_p^0$ ), volume  
 ( $V_p^0$ ), as well as other quantities pertaining to the particular constitutive model used  
 (temperature, plasticity states, etc.). To advance the simulation from time  $t^k$  to time  
 $t^{k+1} = t^k + \Delta t$ , a Galerkin projection of particle momenta to a (usually Cartesian)  
 grid is first carried out, allowing grid velocity to be calculated. This projection is

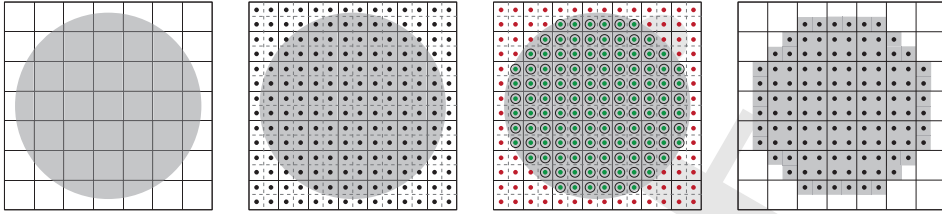


Figure 1: A particle model of a disc is created. (a) The gray area represents the disc, laid on top of the MPM background grid. (b) The grid cells are subdivided into subcells, and the subcell centroids are sampled. (c) The circled samples fall inside the disc and become particles in the model. (d) The particles are shaped like the subcells they were sampled from, so the final model has unavoidably jagged edges.

approximated in MPM as:

$$m_i = \sum_p \phi_i(\mathbf{x}_p) m_p, \quad (1)$$

$$\mathbf{v}_i^k = \frac{\sum_p \phi_i(\mathbf{x}_p^k) \mathbf{v}_p^k m_p}{m_i}, \quad (2)$$

where  $m_i$ ,  $\mathbf{v}_i$ ,  $\phi_i$  are the grid mass, grid velocity, and grid basis functions, respectively. Next, the internal force is computed at the nodes as an approximated volume integral of the divergence of particle stress:

$$\mathbf{f}_i^{\text{int}} = - \sum_p \nabla \phi_i(\mathbf{x}_p^k) \cdot \boldsymbol{\sigma}_p^k V_p^k. \quad (3)$$

External forces (body forces and tractions) are specified on the grid (or projected from particles to the grid), and grid acceleration is computed as:

$$\mathbf{a}_i = \frac{\mathbf{f}_i^{\text{int}} + \mathbf{f}_i^{\text{ext}}}{m_i}. \quad (4)$$

An updated grid velocity is found with a Forward Euler scheme:

$$\mathbf{v}_i^* = \mathbf{v}_i^n + \mathbf{a}_i \Delta t. \quad (5)$$

Next, particle positions and velocities are updated by evaluating the resulting grid velocity and acceleration fields at the particle locations:

$$\mathbf{x}_p^{k+1} = \mathbf{x}_p^k + \sum_i \phi_i(\mathbf{x}_p^k) \mathbf{v}_i^* \Delta t, \quad (6)$$

$$\mathbf{v}_p^{k+1} = \mathbf{x}_p^k + \sum_i \phi_i(\mathbf{x}_p^k) \mathbf{a}_i \Delta t. \quad (7)$$

Evaluation of the gradient of the grid velocity function provides a velocity gradient at particle positions.

$$\nabla \mathbf{v}_p = \sum_i \nabla \phi_i(\mathbf{x}_p^{k+1}) \mathbf{v}_i^*. \quad (8)$$

This velocity gradient is used in updating particle deformation gradient, which is in turn used in calculating particle stress and particle volume:

$$\mathbf{F}_p^{k+1} = (\mathbf{I} + \nabla \mathbf{v}_p \Delta t) \mathbf{F}_p^k, \quad (9)$$

$$\boldsymbol{\sigma}_p^{k+1} = \boldsymbol{\sigma}(\mathbf{F}_p^{k+1}), \quad (10)$$

$$V_p^{k+1} = V_p^0 \det(\mathbf{F}_p^{k+1}), \quad (11)$$

95 where the constitutive model relating deformation to stress is represented by the  
 96 general function  $\boldsymbol{\sigma}$ . Note that while Equation 10 has a dependence only on de-  
 97 formation gradient, implying a hyperelastic material model, the methodology de-  
 98 scribed here is completely general with respect to constitutive model. In other  
 99 words, while it *is* necessary to compute a value for  $\mathbf{F}_p$  as given in Equation 9, any  
 100 constitutive model (e.g., a hypoelastic model with rate dependent plasticity) may  
 101 be used. Even if the deformation gradient is not used to advance the solution, ad-  
 102 vancing it in time is of low cost, and, as the examples below will show, high value.  
 103 In addition, of course, the deformation gradient is valuable in that it can be used to  
 104 compute numerous measures of strain.

105 Equations (1 - 11) outline a single timestep of MPM. GIMP follows the same out-  
 106 line, but the terms  $\phi_i(\mathbf{x}_p)$  are replaced by a different function  $\bar{\phi}_{i_p}$ , calculated as a  
 107 convolution of the grid basis function  $\phi_i$  and a particle characteristic function  $\chi_p$   
 108 (for more details about GIMP and particle characteristic functions, see Barden-  
 109 hagen and Kober (2004)).

110 In the remainder of this paper, we will be working with GIMP simulations, though  
 111 the method we present works with traditional MPM. It will also work with any  
 112 particle method that treats particles as subvolumes of a continua (as opposed to  
 113 sampling points) that is capable of computing a deformation gradient, whether or  
 114 not it is used to advance the solution.

## 115 2 Related Work

116 Bigler, Guilkey, Gribble, Hansen, and Parker (2006) present an overview of their  
 117 methods for visualizing MPM data produced by the Center for the Simulation of

118 Accidental Fires and Explosions project (C-SAFE) [Henderson, McMurtry, Smith,  
119 Voth, Wight, and Pershing (2000)]. They use the SCIRun Problem Solving En-  
120 vironment [Parker and Johnson (1995)] to view smaller data sets using standard  
121 graphics hardware, and a high-performance ray tracer [Bigler, Stephens, and Parker  
122 (2006)] to handle the millions of primitives in larger data sets. In both systems,  
123 each particle is represented by a sphere of a specific color and radius, both chosen  
124 to represent physical parameters in the data set. Generally, in such visualizations  
125 deformation is understood on a global scale, allowing the particle positions to re-  
126 lay an idea of deformation across the whole model, without any indication of the  
127 shapes of individual particles. For understanding smaller structures at closer scales,  
128 lighting and non-photorealistic rendering techniques are employed. Shadowing and  
129 ambient occlusion give visual cues about relative positions of ambiguously oriented  
130 particles, while silhouette edges help to bring attention to medium-scale structures  
131 by highlighting closely bound groups of particles.

132 Gribble, Stephens, Guilkey, and Parker (2006) developed a visualization system  
133 that takes advantage of the symmetry of spheres. They use programmable graphics  
134 hardware to accelerate particle rendering by using a texture mapped billboard of  
135 a single sphere to represent each particle, instead of rendering actual geometry.  
136 By reducing the number of triangles needed to just two per billboard, the system  
137 achieves interactive rates on desktop machines for order-of-million-particle data  
138 sets. The standard method of simply rasterizing millions of triangle-tessellated  
139 spheres quickly overwhelms current graphics hardware; this approach is, therefore,  
140 notable for its high performance. However, as it displays the same prerendered  
141 sphere geometry for every particle, it cannot handle deformation data via glyph-  
142 based tensor visualization as discussed above.

143 In summary, state-of-the-art particle data visualization treats each particle simply  
144 as a position associated with data. This simple approach works well for many  
145 purposes, but this paper will argue that we can build on these methods, including  
146 deformation data in a physically correct and visually insightful way.

147 On the other hand, the graphics community's treatment of deformation tends to  
148 ward approximate but credible images rather than numerical accuracy, since such  
149 approximations are computationally cheaper to achieve. Such methods are used in  
150 settings where accuracy is secondary to visual effect, such as in movies and video  
151 games. Typical methods are based on simulation of continuum and fracture me-  
152 chanics through finite elements for both brittle [O'Brien and Hodgins (1999)] and  
153 ductile [O'Brien, Bargteil, and Hodgins (2002)] fracture. Irving, Teran, and Fed-  
154 kiw (2004) have developed invertible finite elements as an extension to standard  
155 finite elements which behaves robustly enough to handle physical situations with  
156 extreme mesh deformations and even inverted mesh elements, but at a significant

157 cost in the accuracy of the results. These methods are valuable mainly for visual  
 158 results and not scientific accuracy. These and other related methods in graphics can  
 159 be traced back to methods dealing both with elasticity [Terzopoulos, Platt, Barr,  
 160 and Fleischer (1987)] and inelasticity (i.e. viscoelasticity, plasticity, and fracture)  
 161 [Terzopoulos and Fleischer (1988)] in graphical models.

### 162 3 Visualizing Particle Deformation

163 Glyph-based particle visualization methods can be summarized as follows: For  
 164 each particle  $p$  situated at location  $\mathbf{x}_p$ , select a set of points  $G_p$  (representing some  
 165 glyph geometry) centered at the origin, a deformation operator  $\mathbf{D}_p$ , and a color  $C_p$ .  
 166 Render the deformed geometry  $\mathbf{D}_p(G_p)$  translated to location  $\mathbf{x}_p$  with color  $C_p$ .

167 Common particle visualization techniques typically use a unit sphere for  $G_p$ , a  
 168 scaling operation for  $\mathbf{D}_p$  with a magnitude chosen to reflect a scalar data value on  
 169 particle  $p$  (or alternatively, simply no scaling at all), and a scalar color map applied  
 170 to another value from  $p$  to specify  $C_p$ . These choices yield spheres whose radii  
 171 possibly reflect one scalar value (commonly a volume or mass), and whose colors  
 172 reflect another (such as equivalent stress, velocity magnitude, temperature, etc.).

173 In order to visualize deformation, we instead use  $\mathbf{F}_p$ , the deformation gradient of  
 174 particle  $p$ , for the deformation operator by applying it as a linear transformation to  
 175 the points making up the origin-centered glyph. This transformation illustrates the  
 176 local deformation by applying it directly to the piece of the object on which it acts.

177 The key is to select a glyph geometry that communicates information about the  
 178 deformation gradient. For example, spheres do not work for the simple reason that  
 179 they cannot indicate pure rotation. Suppose  $G_S$  is a unit sphere centered at the  
 180 origin, and  $\mathbf{F}$  is a deformation gradient with polar decomposition [Strang (1988)]  
 181  $\mathbf{F} = \mathbf{R}\mathbf{S}$  (in which  $\mathbf{R}$  is orthonormal, representing a rotation, and  $\mathbf{S}$  is symmetric  
 182 and positive-definite, representing non-uniform but orthogonal scaling). Because  
 183  $\mathbf{S}$  is symmetric and positive-definite, it can be decomposed as  $\mathbf{S} = \mathbf{Q}\mathbf{\Lambda}\mathbf{Q}^T$ , where  
 184  $\mathbf{\Lambda}$  is diagonal and  $\mathbf{Q}$  is orthonormal. Because rotational transforms do not affect  
 185 a sphere, we also have  $\mathbf{U}G_S = G_S$  for any orthonormal  $\mathbf{U}$ .<sup>1</sup> Finally, concatenating  
 186 rotation transforms (represented by products of orthonormal matrices) yields a ro-  
 187 tation transform; i.e.  $\mathbf{R}\mathbf{Q}$  can be written as  $\mathbf{U}$  for orthonormal  $\mathbf{R}$  and  $\mathbf{Q}$ , and some

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<sup>1</sup> We have adopted the notation  $\mathbf{F}S$  for the notion of a tensor  $\mathbf{F}$  acting on a set of points  $S$ , which we define as follows:  $\mathbf{F}S \equiv \{\mathbf{F}x | x \in S\}$ .

188 orthonormal  $\mathbf{U}$ :

$$\begin{aligned}
 \mathbf{F}G_S &= \mathbf{R}S\mathbf{G}_S \\
 &= \mathbf{R}\mathbf{Q}\mathbf{A}\mathbf{Q}^T G_S \\
 &= \mathbf{R}\mathbf{Q}\mathbf{A}\mathbf{Q}^T \mathbf{R}^T G_S \\
 &= \mathbf{U}\mathbf{A}\mathbf{U}^T G_S \\
 &= \mathbf{S}'G_S.
 \end{aligned}$$

189 The general transform  $\mathbf{F}$  is aliased by some symmetric positive-definite transform  
 190  $\mathbf{S}'$ . In other words, with respect to spheres, *every linear transformation is a scal-*  
 191 *ing transformation*. If the tensors used for deforming the glyphs are symmetric  
 192 and positive-definite (such as stress), then they lack a rotational component and  
 193 spheres (or other glyphs [Kindlmann (2004)]) can be used illustrate the tensors'  
 194 eigenspaces. For purposes of studying deformation, however, capturing rotation is  
 195 necessary, so a more suitable glyph geometry is required.

196 In addition to properly handling rotation, the glyph geometry should also repre-  
 197 sent the initial discretization behind the structural mechanics simulation. Often  
 198 spheres are used to visualize particle data because particles often represent geo-  
 199 metric points. The word “point” in “material point method” emphasizes this sim-  
 200 plifying idea. However, MPM models continuum mechanics, and so rather than  
 201 being dimensionless points, MPM particles represent computational volumes that  
 202 initially partition a continuous object.

203 The common choice of sphere geometry therefore has geometric and representa-  
 204 tional deficiencies. By examining how MPM often initializes the particles in a  
 205 particle model, we conclude that *cuboid* glyphs produce meaningful and insight-  
 206 ful visualizations. Furthermore, we take advantage of the generality inherent in  
 207 MPM’s modeling algorithm to create *hexahedral* particle models that do not suffer  
 208 from the axis-alignment constraint placed upon cuboid particle models.

### 209 3.1 Cuboid Glyphs

210 When initializing a particle model, GIMP requires that each particle have a position  
 211 and a volume (which is used to normalize integrals over other particle values).  
 212 GIMP has an implicit notion of the shapes of particles (as described below, and in  
 213 Figure 1), and we argue it is useful for visualization, in addition to computation.

214 A standard way to create an MPM model regularly divides the background grid  
 215 cells into subcells (Figure 1). If the centroid of a subcell falls inside the boundary  
 216 of the object being modeled, then a particle is initialized at the centroid location,  
 217 with volume equal to the volume of the subcell. Because the initial volume is



218 derived from the subcell's shape, the modeling process further implies that each  
219 particle actually looks like a cuboid with the same dimensions as the subcell. In  
220 other words, a *cuboid* glyph geometry is suitable for visualizing this model. Using  
221 cuboids in this case produces a visually continuous model occupying exactly the  
222 volume the modeling process dictates. The particles now look like the continua of  
223 matter they represent, connected at the faces in such a way that the particle glyphs  
224 reflect visually the manner in which the object is partitioned numerically.

225 Cuboids are used primarily because they reflect the underlying numerical represen-  
226 tation of the material, but they also have desirable graphical properties that enhance  
227 the visualization. Cuboids are bounded by quadrilateral faces, enabling more effec-  
228 tive visualization of deforming surfaces. For instance, the exposed faces of cuboids  
229 on the outer edge of an object represent its surface; as the object deforms, lighting  
230 effects help the viewer visually track the surface as it changes shape. Put another  
231 way, sphere glyphs allow a viewer to track a deforming surface, but only by their  
232 relative positions; cuboids, on the other hand, show relative positions *and* relevant  
233 lighting cues to give a much stronger sense of a deforming surface. Furthermore,  
234 in a cuboid glyph model, the edges of the cuboids form a kind of grid of junction  
235 lines that is visible in visualizations. As the simulation progresses, these junction  
236 lines change their shape to reflect the changing shape of the particles (Figure 2);  
237 they can also be used to track deformation on a larger scale (Figure 3).

### 238 3.2 Hexahedral Glyphs

239 Because the MPM background mesh is most often a rectilinear grid, the bound-  
240 ary of the particle models produced by the standard modeling algorithm are also  
241 constrained to be rectilinear. Models made to approximate objects with curved  
242 boundaries will have a stairstep or lego-brick quality along these curves (Figure 1).



Figure 2: A cylinder, discretized in CUBIT, is compressed by rigid plates (not shown) until it buckles, introducing a large deformation.

243 Initially, the particle glyphs will be axis-aligned; instead of seeing the surface of a  
244 sphere, for example, we will only see a stairstep approximation to that surface. By  
245 generalizing cuboids to *hexahedra*, however, we can remove this restriction on the  
246 initial orientation of the particle glyphs.

247 CUBIT [Sandia National Laboratories (2007)] is a software package that can pro-  
248 duce hexahedral meshes.<sup>2</sup> Such meshes have several desirable properties for MPM  
249 particle models: they approximate the interior of some boundary, covering it con-  
250 tinuously with volumetric elements, and the exposed faces of the boundary mesh  
251 elements can be made to approximate a chosen surface. To create a particle model  
252 from a hex mesh, we simply initialize

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<sup>2</sup> Of the ten topological classes of hexahedra, this paper uses “hexahedron” to mean “quadrilateral-faced hexahedra,” the variety topologically equivalent to a cube.

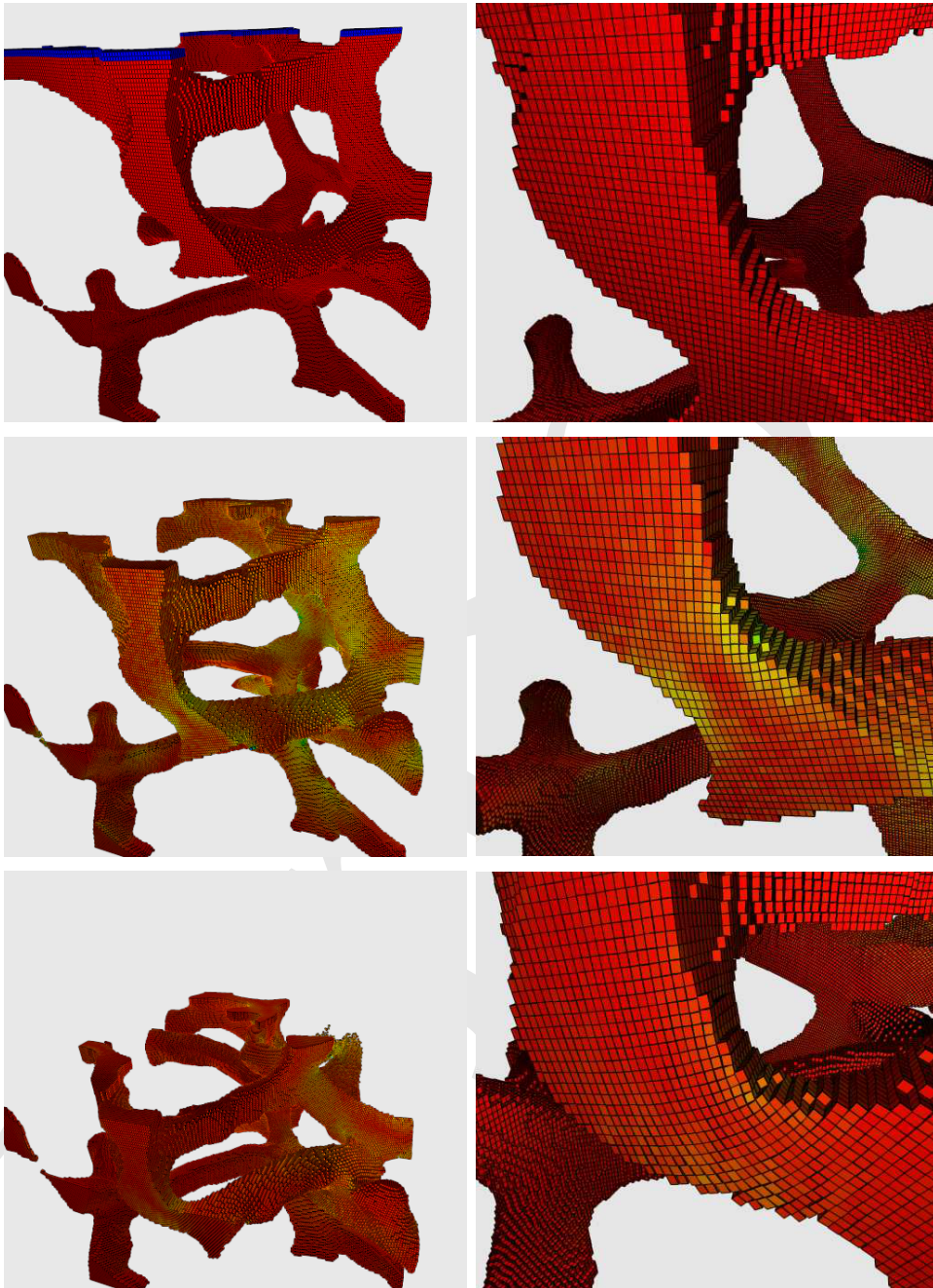


Figure 3: Open-celled foam made up of microstruts is compacted by a rigid plate moving downward (not shown). The struts deform in bearing the load as the total volume of the foam decreases. The strut in the left of the foreground deforms considerably during the process (detail, right column). This strut is originally vertical and bends during compression. In the lower right image the rotation associated with this deformation can be observed by tracking the boundaries between particles. Along the bottom edge of the strut, where tensile stresses are acting, some particle separation can be seen, possibly reflecting some numerical inaccuracies in the simulation.

253 one particle at the geometric centroid of each mesh element, with its initial volume  
254 set to that of the mesh element. This list of material points is given to the MPM  
255 algorithm, which can then carry out a simulation. CUBIT can create meshes to rep-  
256 resent cuboids, cylinders, prisms, cones, polygonal pyramids, spheres, and toruses,  
257 and it can compose these primitives into more complex shapes using constructive  
258 solid geometry. It can therefore offer a very general range of particle models. To  
259 perform the visualization, each particle is represented by a hexahedral glyph with  
260 the same shape as the mesh element that produced it and then transformed about its  
261 centroid by the deformation gradient tensor. This approach allows the simulation  
262 scientist to create more realistic-looking models, which may lead to better insight  
263 gained from visualization.

## 264 4 Examples and Discussion

265 To test our method, we have run several MPM simulations with different geometry  
266 and physical conditions. Some of the simulations aim to produce specific modes  
267 of deformation for observation, while others are real data produced by the C-SAFE  
268 project [Henderson, McMurtry, Smith, Voth, Wight, and Pershing (2000)]. The  
269 images were all produced by the Manta ray tracer [Bigler, Stephens, and Parker  
270 (2006)], which includes spheres, cuboids, and hexahedra as graphical primitives,  
271 and runs at interactive frame rates on modest desktop hardware. We used a variation  
272 of the standard ray tracing algorithm to also render non-photorealistic intersection,  
273 crease, and silhouette lines [Choudhury and Parker (2009)]; these lines help to show  
274 the spatial relationships of the individual glyphs.

275 Figure 3 gives an example of our method applied to real data. The simulation shown  
276 involves a small volume of foam, whose microstruts are visible, being crushed by  
277 a downward-moving rigid plate. The geometry for the foam model was created by  
278 obtaining X-ray microtomography data of a real foam sample, then using imaging  
279 techniques on the resulting data volume and initializing particles for voxels surpass-  
280 ing a threshold intensity (indicating the presence of foam in that voxel) [Brydon,  
281 Bardenhagen, Miller, and Seidler (2005)]. As the particles in this case were de-  
282 rived from voxel data, we use cuboid glyphs for visualization. The right column  
283 of Figure 3 shows a close-up of one strut that deforms considerably during the  
284 simulation.

### 285 4.1 Physical Basis

286 The primary feature of our visualizations is that they are physically based, deter-  
287 mining glyph shape from the modeling process, and then deforming the glyphs in  
288 accordance with deformation data generated during the simulation. In particular,  
289 the method provides a *volumetric view* of the particles, showing how they would

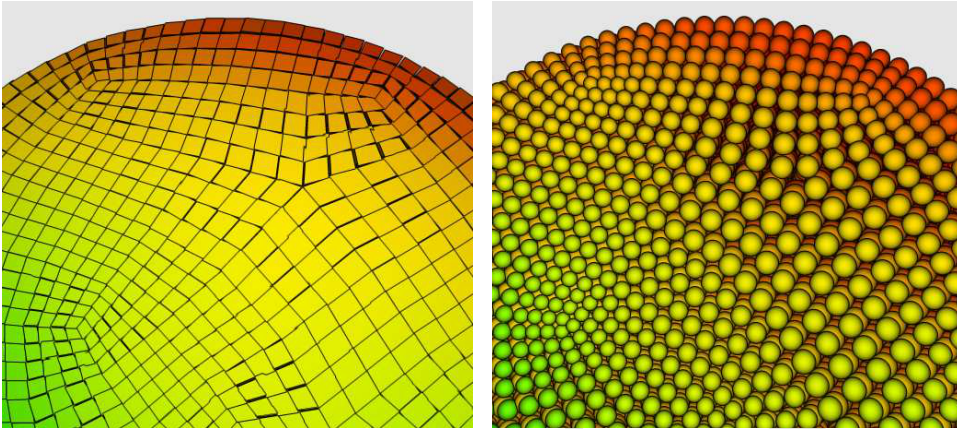


Figure 4: Bottom of cylinder depicted in Figure 2. Hexahedral glyphs emphasize the volumetric nature of MPM particles; spheres emphasize their pointwise nature. In particular, slight particle separation is visible in the left image; such separations are not apparent in the right image.

290 appear under the assumption that the entire particle voxel deforms with a constant  
 291 deformation gradient, as computed at the particle's nominal position. Figure 4  
 292 shows the bottom surface of a cylinder that is being being crushed longitudinally  
 293 by two rigid plates (as shown in Figure 2). The particles in this part of the cylin-  
 294 der are experiencing loads that tend to flatten them out. The hexahedral glyphs in  
 295 Figure 4, left, are therefore flat and thin, so that they occupy more screen space  
 296 than their spherical counterparts in the right image. On the other hand, in Figure 4,  
 297 right, one can actually see past the spheres into the interior of the model, consti-  
 298 tuting a *computational view* of the particles that emphasizes their pointwise nature,  
 299 as they exist during computation by the MPM algorithm. By using small, volume-  
 300 independent spherical glyphs, we can observe the relative positions of particles  
 301 throughout the volume of a simulated body (for this purpose, global illumination  
 302 can also help [Gribble, Stephens, Guilkey, and Parker (2006); Gribble and Parker  
 303 (2006); Tarini, Cignoni, and Montani (2006)]).

304 The major strength of hexahedral glyphs is that they give a good physical picture of  
 305 what the simulated object might look like in reality, which also has other implica-  
 306 tions. For example, MPM simulates continuum bodies, which means that material  
 307 separation does not occur unless a material failure model is included. The parti-  
 308 cles in MPM are not connected as in a Finite Element Mesh, thus, the degree to  
 309 which edges of cuboid particles remain connected visually is a reflection that the  
 310 method is accurately capturing the deformation that an object experiences. How-

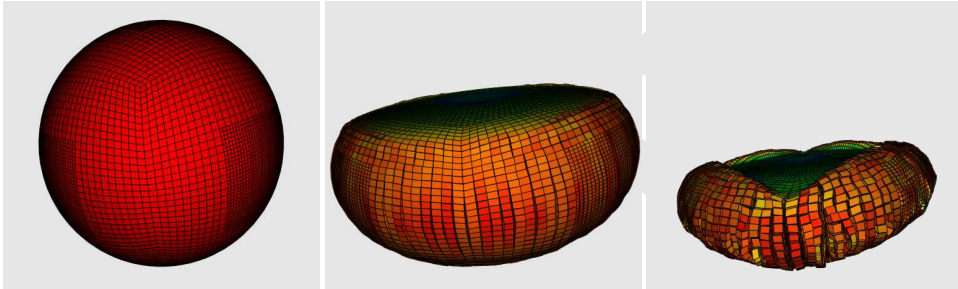


Figure 5: A sphere is compressed by rigid plates (not shown) until it is flat in the middle. Some material oozes out from between the plates, spreading out in such a way that the particles become separated. Because the material model did not include a failure model in this simulation, the particle separation indicates the approximate nature of the simulation. The larger separations in the rightmost image suggests possible error as well.

311 ever, as the simulation is necessarily an approximation to the true behavior, we may  
 312 observe separation between particles in the resulting visualization. Such visual sep-  
 313 aration reflects both the approximate nature of the simulation results, as well as the  
 314 visualization assumption that the whole particle deforms according to a constant  
 315 per-particle value of  $\mathbf{F}$  (Figure 4, left). However, large, localized separations can  
 316 indicate error in the simulation (Figure 5).

317 Such separations are not apparent when using spherical glyphs because spheres do  
 318 not properly model the material continuum. This is a case of the visualization draw-  
 319 ing attention to errors in the simulation, providing hints to the simulation scientist  
 320 about how the simulation quality might be improved. In Figure 5, right, the sepa-  
 321 ration is quite large and may, for example, indicate the need to run the simulation  
 322 again for higher accuracy, using a discretization with a larger number of smaller  
 323 particles in that area.

324 Several of the simulations demonstrate the usefulness of using a volumetric view.  
 325 Figure 6 shows the results of a Taylor impact [Taylor (1948)] simulation, in which  
 326 a metal cylinder is shot onto a rigid surface and deforms. Because this setup is  
 327 radially symmetric about the cylinder's axis, the simulation uses only one quarter  
 328 of the cylinder with appropriate boundary conditions. Looking at the internal corner  
 329 of the model actually shows what is happening in the center of the cylinder. The  
 330 extremely high compressive stress in this region induces the particles to become  
 331 very wide and thin. In the visualization, the particles remain continuous, suggesting  
 332 the numerical stability of the simulation.

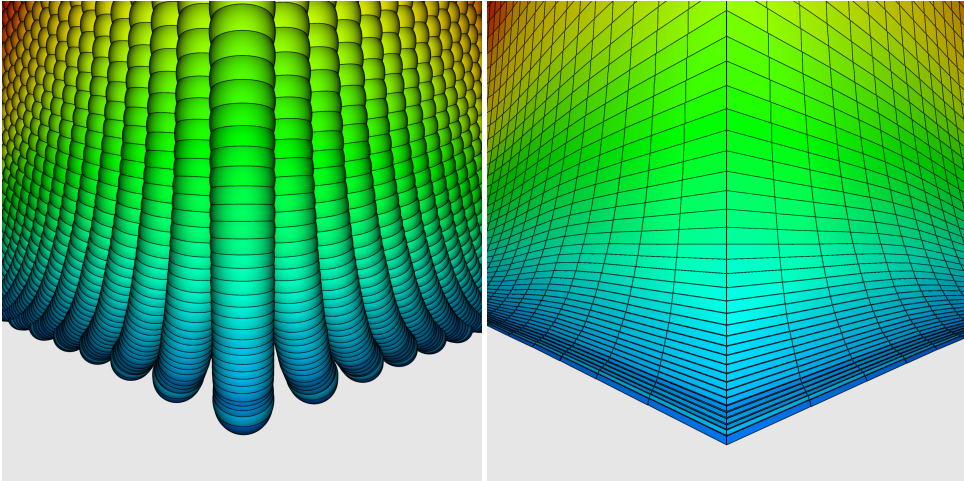


Figure 6: Extreme deformation is induced by very high compressive stresses in the center of a Taylor-impacted cylinder. By changing their shape to be flat and wide, cuboid glyphs are able to illustrate the deformation; by comparison, the sphere glyphs become impacted, hiding each other from view, and imply that the columns of particles are becoming separated.

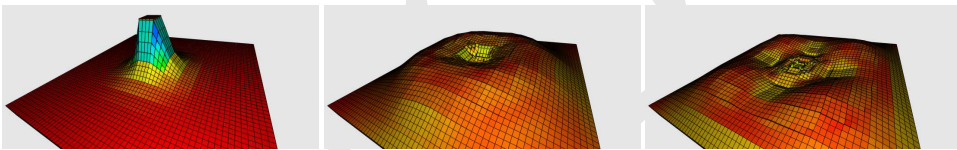


Figure 7: A rubber sheet fixed around its edges to a wall, struck by a projectile from underneath. An elastic wave travels outward from the point of impact.

333 Figure 7 shows a rubber sheet with its edges affixed to a wall being struck by a  
 334 projectile from underneath. The middle part of the sheet moves upward with the  
 335 projectile, and then returns to its original position while an elastic wave travels  
 336 out from the point of impact. The graphical qualities of the cuboid glyphs nicely  
 337 illustrate the curve of the wave as it travels. The wave travels outward radially, but  
 338 the front strikes different parts of the boundary at different times, taking longer to  
 339 reach the corners of the sheet than the sides. The reflected waves therefore return at  
 340 different times and cause a characteristic interference pattern, as illustrated by the  
 341 shading in the rightmost image of Figure 7.

342 Furthermore, this particular visualization reveals something interesting about the  
 343 MPM algorithm itself. It is subtly apparent in Figure 7 that the particles are some-

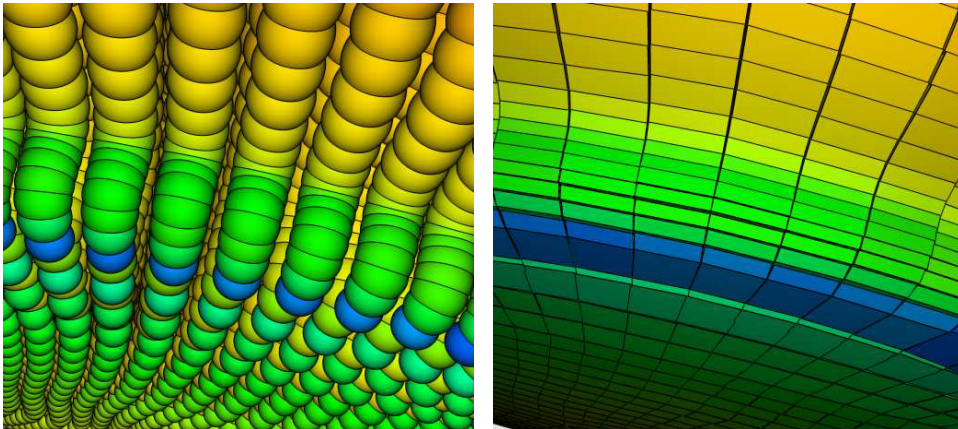


Figure 8: The spheres falsely imply that the corner of the cylinder lies along the green particles, while the hexahedra show that actually, the corner occurs in the blue particles, and the green particles show a slight bulge just above it.

344 how organized into “tiles” with similar surface orientation. These are visible from  
 345 the changes in shading resulting from the slight change in the angles of the tiles  
 346 with respect to each other. The effect results from the use of the background grid  
 347 used in MPM, which gathers and then reprojects changes to particle state through  
 348 basis functions during the simulated timestep. This is an example of the visual-  
 349 ization revealing a quality of the algorithm itself, in much the same way as small  
 350 gaps appearing between particles reflect the approximating nature of any simulation  
 351 algorithm.

#### 352 **4.2 Visual Cues for Geometric Features**

353 When studying deformation in simulation data, it is important that the visualiza-  
 354 tion not misrepresent or obscure important features. Hexahedral glyphs are a way  
 355 to eliminate such misrepresentations, as demonstrated in the bottom corner of the  
 356 cylinder (Figure 8). In the sphere scene, the cylinder’s corner seems to occur along  
 357 the row of green particles three layers above the blue ones, but the hex scene  
 358 demonstrates that this is not the case. In fact, the corner appears in the row of  
 359 blue particles, as evidenced by the two faces visible in that row. What looks like  
 360 the corner in the sphere scene is in fact a bulge that occurs just above the plane of  
 361 contact with the rigid plate. The shading and orientation of the hex faces shows this  
 362 phenomenon quite clearly.

363 When viewing a dataset at close ranges to investigate small-scale features, global  
 364 indicators of structure and deformation are missing. The inside of the buckle feature



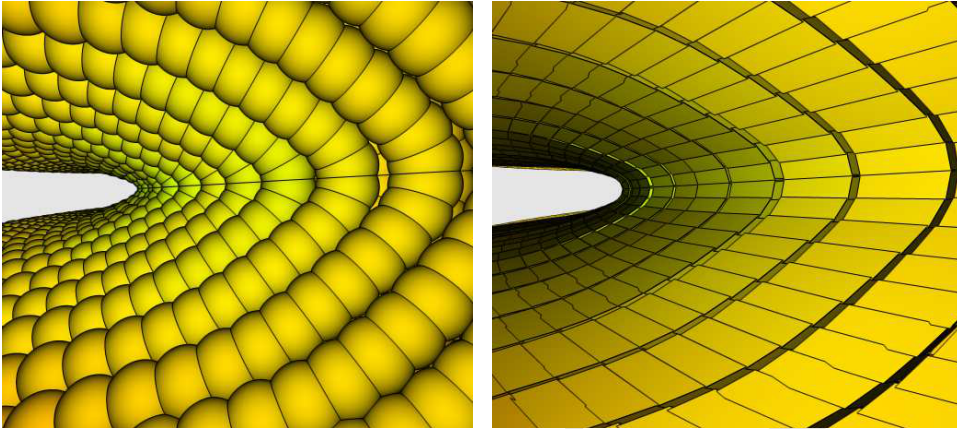


Figure 9: Hexahedra show the geometry of high curvature areas more clearly than spheres can through shading cues that suggest surfaces directly.

365 in the crushed cylinder (Figure 9) demonstrates this problem. A sphere's surface  
 366 normal varies continuously across its surface; many spheres packed close together  
 367 in a visualization will therefore look very similar. For this reason, in the left image  
 368 of Figure 9 it is very difficult to infer the actual distribution of the glyphs. It is  
 369 not clear if the curved bands of connected spheres lie in the same plane, or if they  
 370 recede from the viewpoint. The structure is much easier to see in the right image,  
 371 because the exposed faces of the hexahedral glyphs have one normal vector each,  
 372 and thus can serve as area elements making up a surface, even at close viewing  
 373 distances.

## 374 5 Conclusions and Future Work

375 We have demonstrated an extension to glyph-based particle visualization meth-  
 376 ods that includes the deformation gradient and therefore visualizes deformation  
 377 directly. Our major strategy has been to understand a given particle method's ap-  
 378 proach to modeling the geometry of an individual particle, and then adopt that ge-  
 379 ometry as the visualization glyph. This strategy moves particle visualization away  
 380 from arbitrary choices, such as spheres, that seem to be appropriate, but turn out to  
 381 be deficient upon close inspection.

382 By visualizing the particles as the hexahedral regions they represent rather than  
 383 abstract points in space, scientists can get a clearer and more direct understanding  
 384 of how a simulation affects the material being modeled. This directness is espe-  
 385 cially important when deformation is the central object of interest, as in the study

386 of material failure during a catastrophic explosion, or stress testing of structures.  
387 Furthermore, including deformation data directly in the visualization allows simu-  
388 lation scientists to detect errors in their simulations earlier, as when the visualiza-  
389 tion shows deformations that clearly do not reflect reality.

390 In future work we wish to integrate other common visualization techniques. For  
391 instance, it has been demonstrated that using ambient occlusion or other approx-  
392 imations to global illumination can enhance user perception of particle datasets  
393 [Gribble and Parker (2006); Tarini, Cignoni, and Montani (2006)]; presumably this  
394 is true when using hexahedra rather than spheres as well.

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