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#### Abstract

An analytical approach to solving anisotropic single scattering from point light sources in homogeneous media was recently derived via a dual-formulation of the air-light integral. In this paper, we demonstrate how to reduce the evaluation of the terms involved in the solution and provide an efficient and practical implementation substantially increasing the real-time performance characteristics.

CR Categories: I.3.7 [Computer Graphics]: Three-Dimensional Graphics and Realism-Color, shading, shadowing, and texture Keywords: participating media, analytical integration


## 1 Introduction \& Related Work

[Pegoraro and Parker 2009] recently derived a closed-form solution to single scattering in homogeneous isotropic media. They subsequently extended the analytical approach to anisotropic phase functions and light distributions via domain partitioning of the air-light integral [Pegoraro et al. 2009]. Solving the resulting dual-formulation however requires multiple evaluations of the computationally costly complex-valued exponential integral function.

In this paper, we reformulate the terms involved in solving the dual-formulation and reduce the computation to only 3 evaluations of the complex-valued exponential integral instead of 4 . In addition, we provide an efficient implementation yielding a substantial speed-up of the real-time performance achieved on graphics hardware.

## 2 Reduced Evaluation \& Implementation

Defining $u$ as the variable of integration in the simplified form of the air-light integral as illustrated in figure 1, [Pegoraro et al. 2009] proposed the changes of variable $v=u+\sqrt{1+u^{2}}$ and $w=u-\sqrt{1+u^{2}}$, and formulated a solution requiring 4 evaluations of the complex-valued exponential integral Ei with the parameters $v_{a}, v_{h}, w_{b}$ and $w_{h}$ involved by the computation of the following terms

$$
\begin{align*}
& I_{0}\left(-H, v_{h}, v_{a}\right)=i_{0}\left(-H, v_{a}\right)-i_{0}\left(-H, v_{h}\right) \\
& I_{1}\left(-H, v_{h}, v_{a}\right)=i_{1}\left(-H, v_{a}\right)-i_{1}\left(-H, v_{h}\right) \\
& J_{0}\left(H, w_{h}, w_{b}\right)=j_{0}\left(H, w_{b}\right)-j_{0}\left(H, w_{h}\right) \\
& J_{1}\left(H, w_{h}, w_{b}\right)=j_{1}\left(H, w_{b}\right)-j_{1}\left(H, w_{h}\right)  \tag{4}\\
& J_{e}\left(H, w_{h}, w_{b}\right)=\operatorname{Ei}\left(\frac{H}{w_{b}}\right)-\operatorname{Ei}\left(\frac{H}{w_{h}}\right) \tag{5}
\end{align*}
$$

with $H$ the optical distance from the light to the ray and where

$$
\begin{aligned}
& i_{0}(a, v)=\sin (a) \Re(\operatorname{Ei}(a v+\imath a))-\cos (a) \Im(\operatorname{Ei}(a v+\imath a)) \\
& i_{1}(a, v)=\cos (a) \Re(\operatorname{Ei}(a v+\imath a))+\sin (a) \Im(\operatorname{Ei}(a v+\imath a)) \\
& j_{0}(a, w)=-\sin (a) \Re\left(\operatorname{Ei}\left(\frac{a}{w}+\imath a\right)\right)+\cos (a) \Im\left(\operatorname{Ei}\left(\frac{a}{w}+\imath a\right)\right) \\
& j_{1}(a, w)=\cos (a) \Re\left(\operatorname{Ei}\left(\frac{a}{w}+\imath a\right)\right)+\sin (a) \Im\left(\operatorname{Ei}\left(\frac{a}{w}+\imath a\right)\right)-\operatorname{Ei}\left(\frac{a}{w}\right) .
\end{aligned}
$$

We here highlight that $1 / w=-v$, that $\cos (-a)=\cos (a)$ and $\sin (-a)=-\sin (a)$, and that $\operatorname{Ei}(\bar{z})=\overline{\operatorname{Ei}(z)}$ yielding the identities

$$
\begin{align*}
& j_{0}(a, w)=i_{0}(-a, v)  \tag{6}\\
& j_{1}(a, w)=i_{1}(-a, v)-\operatorname{Ei}(-a v)
\end{align*}
$$

from which directly follows that the terms involving the complex-valued exponential integral at $v_{h}$ and $w_{h}$ are identical. This reduces the computation to only 3 evaluations of the latter function as illustrated in figure 2 which additionally explicitly sets $u_{h}=0$.


Figure 1: Illustration of the terms involved in the computation of the air-light integral.

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ComputeIJ(v, w
```

1. $c o=\cos (-H)$;
2. $s i=\sin (-H)$;
3. $\{r e, i m\}=\operatorname{Ei}\left(-H v_{a^{\prime}}, H\right)$;
4. $I_{0}=s i * r e-c o *{ }^{\frac{a}{i}}$;
5. $I_{1}=c o * r e+s i * i m ;$
6. $\{r e, i m\}=\operatorname{Ei}\left(H / w_{b}, H\right)$;
7. $J_{0}=s i * r e+c o * i m ;$
8. $J_{1}=c o * r e-s i * i m ;$
9. $\{r e, i m\}=\operatorname{Ei}(-H,-H)$;
10. $I_{0}-=s i * r e-c o * i m ;$
11. $J_{0}-=s i * r e-c o * i m ;$
12. $I_{1}-=c o * r e+s i * i m ;$
13. $J_{1}-=c o * r e+s i * i m ;$
14. $J_{e}=\mathrm{E}_{1}(H)-\mathrm{E}_{1}\left(-H / w_{b}\right)$;
15. $J_{1}-=J_{e}$;
16. return $\left\{I_{0}, I_{1}, J_{0}, J_{1}, J_{e}\right\}$;

Figure 2: Pseudo-code of the reduced evaluation.

## 3 Results \& Conclusion

The method was implemented in a fragment shader using OpenGL and Cg running on an NVIDIA GeForce GTX 280 under Windows Vista 64-bit. The integral was evaluated independently for each color channel, hence 3 times per fragment at a resolution of $512 \times 512$ as illustrated in figure 3 . While the previous method renders at 63 FPS, our new evaluation scheme achieves a frame rate of 77 FPS and consequently yields a speed-up of 1.22X.

In conclusion, we have shown how to reduce the evaluations involved in analytically solving the dual-formulation of the air-light integral for anisotropic single scattering from point light sources in homogeneous media. Moreover, we have provided a practical implementation and demonstrated substantial gains in performance


Figure 3: A lighthouse in thick brume with an anisotropic two-lobed spotlight rendered in real-time.

## References

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