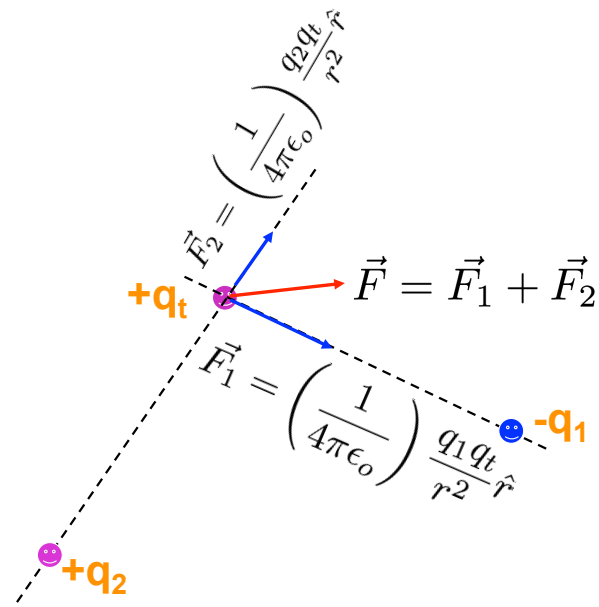


Fields



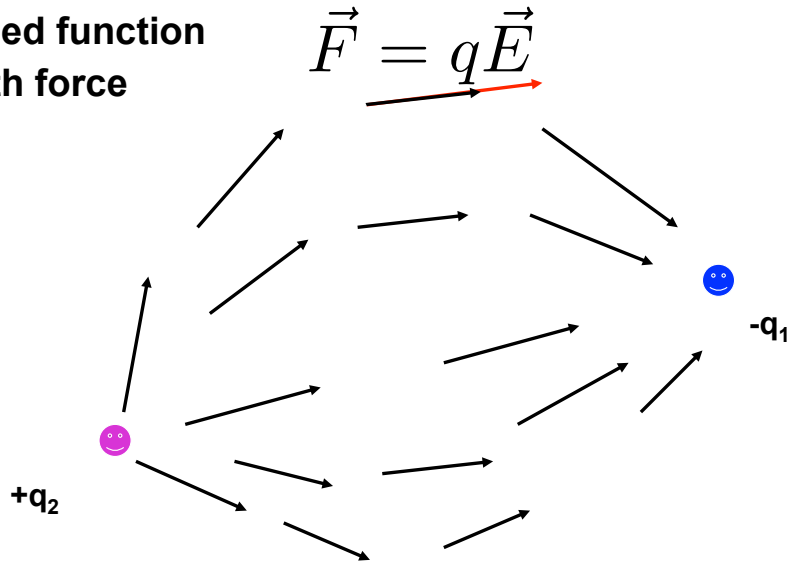
Coulomb's Law



Electric Field

- Vector valued function
- Aligned with force

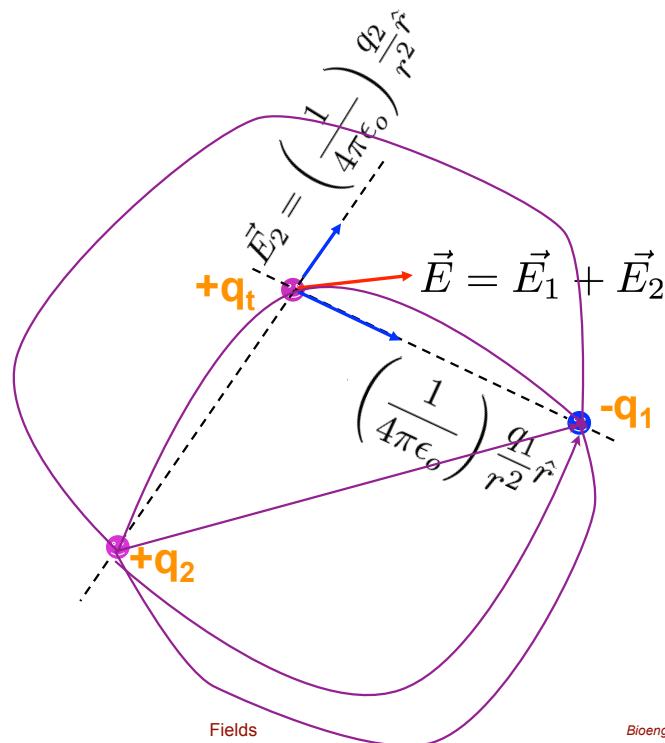
$$\vec{F} = q\vec{E}$$



Fields

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Superposition of Electric Field



Fields

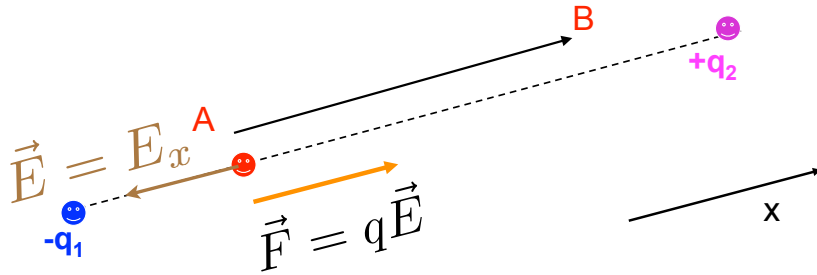
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Potential Energy

$$\Delta U = U(B) - U(A) = -q \int_A^B E_x dx$$

$$dU = F_x dx = qE_x dx$$

$$\Delta U = F_x \Delta x = qE_x \Delta x$$



Fields

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Electric Potential

- Scalar field
- Generalization of potential energy as U is a generalization of F
- Always relative measure
- Note sign convention

Potential Energy

$$\Delta U = U(B) - U(A) = -q \int_A^B E_x dx$$

Electric Potential

$$\Delta u = - \int_A^B E_x dx$$

Another definition of Electric Field

$$E_x = - \frac{\partial u}{\partial x}$$

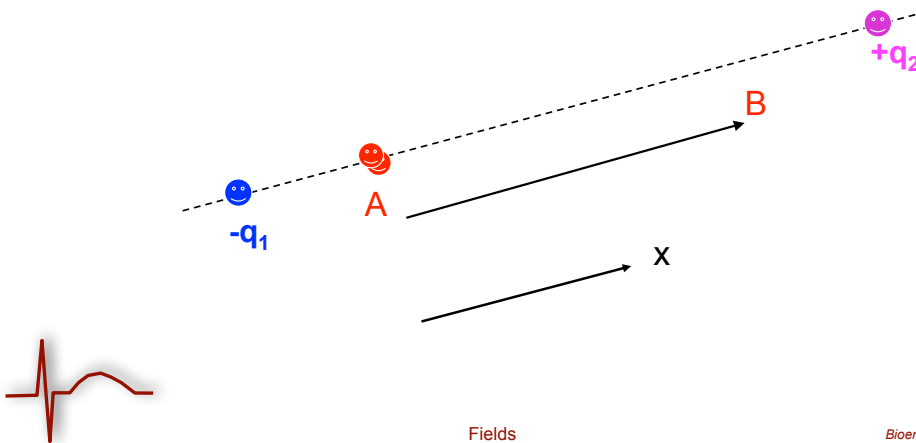


Fields

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Electric Potential Field

- Scalar valued field
- Related to energy and work
- Independent of path



Ohm's Law in a Volume Conductor

- Continuous form of discrete version
- Expressed in terms of field parameters

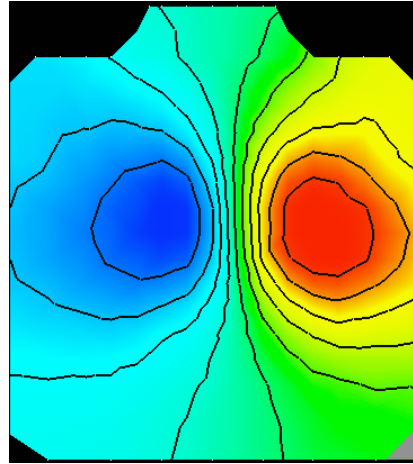
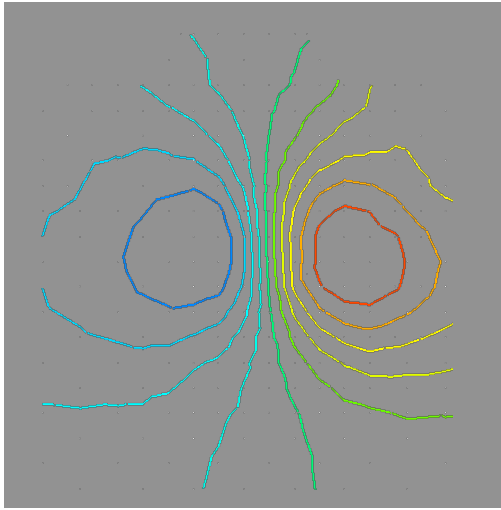
$$j = \frac{i}{A} = \frac{v(B) - v(A)}{RA}$$

$$j_x = -\sigma \frac{\partial v}{\partial x}$$

$$\vec{j} = -\sigma \vec{E}$$



Scalar Fields and Isocontours

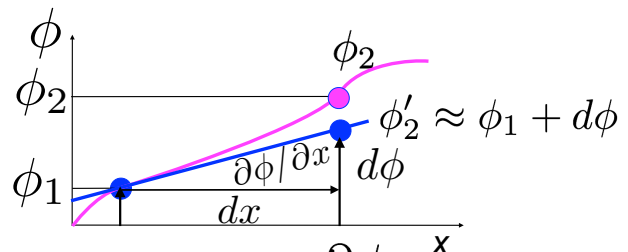
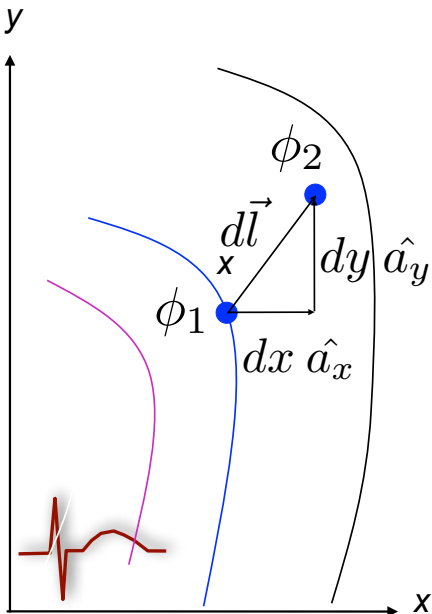


Each line represents a region of constant scalar value

Fields

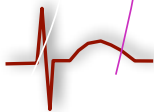
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Scalar Difference



$$\phi'_2 = \phi_1 + \frac{\partial \phi}{\partial x} dx (+ \dots)$$

$$\phi_2 = \phi_1 + \frac{\partial \phi}{\partial x} dx + \frac{\partial \phi}{\partial y} dy (+ \dots)$$



Fields

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The Gradient

We wrote for the scalar difference:

$$\begin{aligned}\phi_2 - \phi_1 &= d\phi \\ &= \partial\phi/\partial x dx + \partial\phi/\partial y dy\end{aligned}$$

Now with

$$d\vec{l} = dx \vec{a}_x + dy \vec{a}_y$$

we can generalize this as a scalar product by defining the vector operator:

$$\nabla\phi = \partial\phi/\partial x \vec{a}_x + \partial\phi/\partial y \vec{a}_y$$

So that the product becomes

$$d\phi = \nabla\phi \cdot d\vec{l}$$

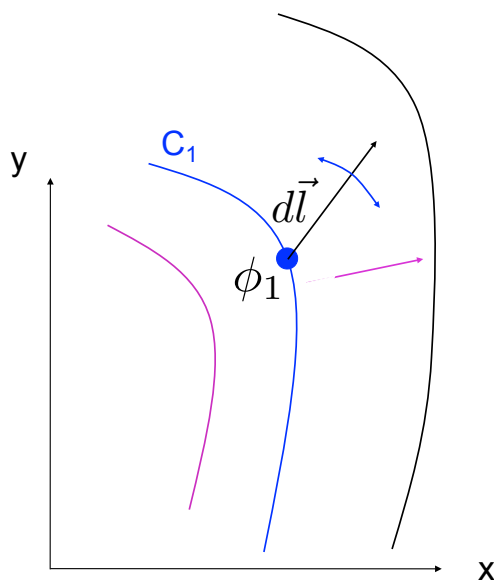
We refer to $\nabla\phi$ as the gradient of the scalar field ϕ . (It is another vector field, i.e., associated with all points in space.)



Fields

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Features of the gradient



$$d\phi = \nabla\phi \cdot d\vec{l}$$

Along which $d\vec{l}$ is $d\phi$ a maximum?
- a minimum?

What is the direction of $\nabla\phi$ relative to the iso-value contour?

If ϕ is electric potential, what is $-\nabla\phi$?



Fields

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Computing the Gradient

$$\nabla \phi = \frac{\partial \phi}{\partial x} \vec{a}_x + \frac{\partial \phi}{\partial y} \vec{a}_y$$

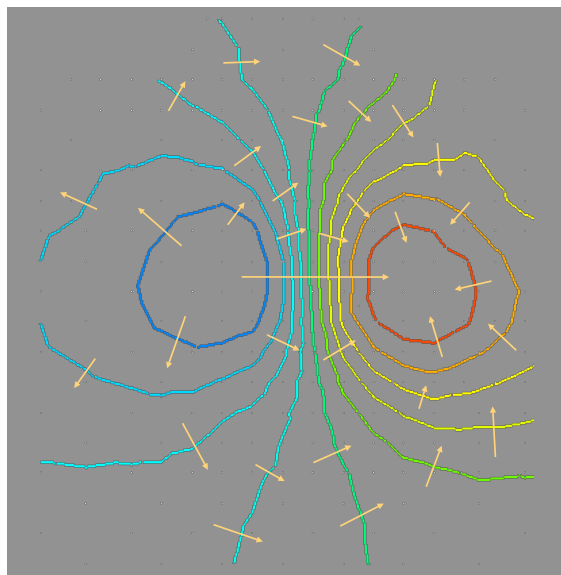
Only in Cartesian coordinates!



Fields

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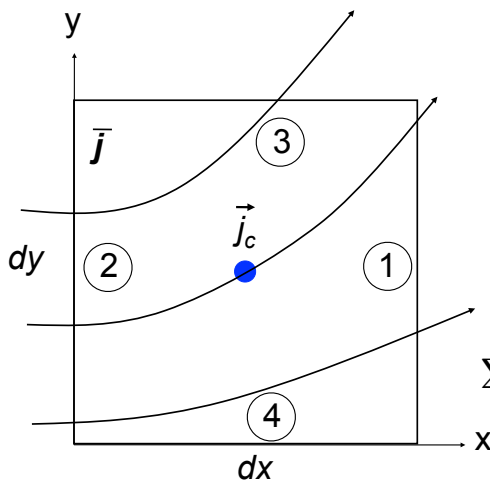
Gradient Example



Fields

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Divergence



$$\text{Flux}_1 = (j_{cx} + \frac{1}{2} \frac{\partial j}{\partial x} dy + \dots)$$

$$\text{Flux}_2 = -(j_{cx} - \frac{1}{2} \frac{\partial j}{\partial x} dy + \dots)$$

$$\text{Flux}_3 = (j_{cy} + \frac{1}{2} \frac{\partial j_y}{\partial y} dx + \dots)$$

$$\text{Flux}_4 = -(j_{cy} - \frac{1}{2} \frac{\partial j_y}{\partial y} dx + \dots)$$

$$\Sigma \text{Flux} = (\partial j_x / \partial x + \partial j_y / \partial y) dx dy$$

But the value of flux depends on $dx dy$, so we define a normalized quantity:

$$\text{Div } \vec{j} = \lim_{xy \rightarrow 0} \frac{\Sigma \text{Flux}}{dx dy} = \frac{\partial j_x}{\partial x} + \frac{\partial j_y}{\partial y} = \nabla \cdot \vec{j}$$



Fields

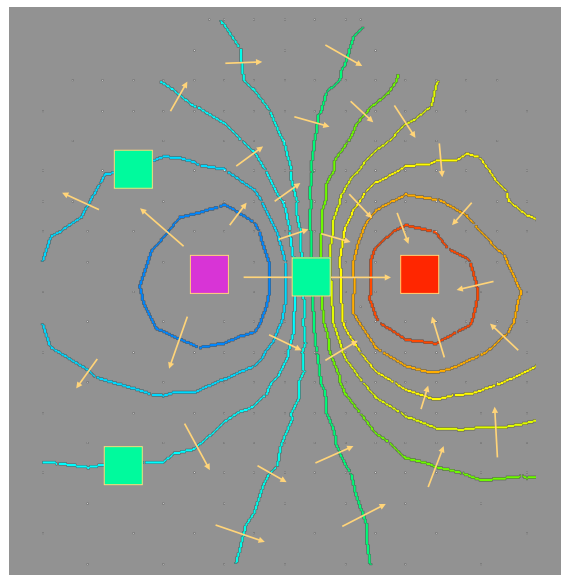
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Divergence Example

■ Div = 0

■ Div > 0

■ Div < 0



Fields

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Gradient to Laplacian

Gradient
$$\vec{\nabla}\phi = \frac{\partial\phi}{\partial x}\vec{a}_x + \frac{\partial\phi}{\partial y}\vec{a}_y$$

Divergence
$$\nabla \cdot \vec{j} = \frac{\partial j_x}{\partial x} + \frac{\partial j_y}{\partial y}$$

Divergence of the gradient = Laplacian
$$\nabla \cdot \vec{\nabla}\phi = \frac{\partial^2\phi}{\partial x^2} + \frac{\partial^2\phi}{\partial y^2} = \nabla^2\phi$$

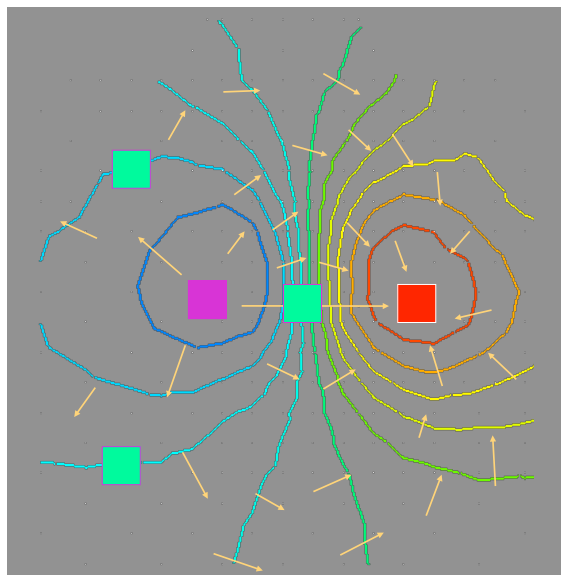


Fields

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Example of Laplacian

- Lap = 0
- Lap > 0
- Lap < 0



Fields

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Back to (Bioelectricity) Basics

- Apply the vector calculus to currents, electric potentials
 - extend to three dimensions
 - note: formulas apply only to Cartesian coordinates
- Bioelectric volume conductors
 - for there to be current, there must be potential difference
 - quasi static assumption (impedance = resistance)
 - may be uniform, piecewise uniform, or anisotropic
 - “regular” circuits are discrete/lumped version of volume conductors



Fields

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Electric Potential

- Scalar field
- Generalization of potential energy as E is a generalization of F
- Always relative measure
- Note sign convention

Potential Energy

$$\Delta U = U(B) - U(A) = -q \int_A^B E_x dx$$

Electric Potential

$$\Delta u = - \int_A^B E_x dx$$

Another definition of Electric Field

$$E_x = - \frac{\partial u}{\partial x}$$



Fields

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Potentials, Fields, Currents in Volume Conductors

$$E_x = -\frac{\partial \phi}{\partial x} \quad \vec{E} = -\nabla \phi \quad \text{Definition of } E$$

$$\vec{J} = \sigma \vec{E} = -\sigma \nabla \phi \quad \text{Ohm's Law}$$

$$\nabla \cdot \vec{J} = -\nabla \cdot \sigma \nabla \phi = -\sigma \nabla^2 \phi = I_v \quad I_v \text{ is source density}$$

$$\nabla^2 \phi = -\frac{I_v}{\sigma} \quad \text{Poisson's Equation}$$

$$\nabla^2 \phi = 0 \quad \text{Laplace's Equation}$$

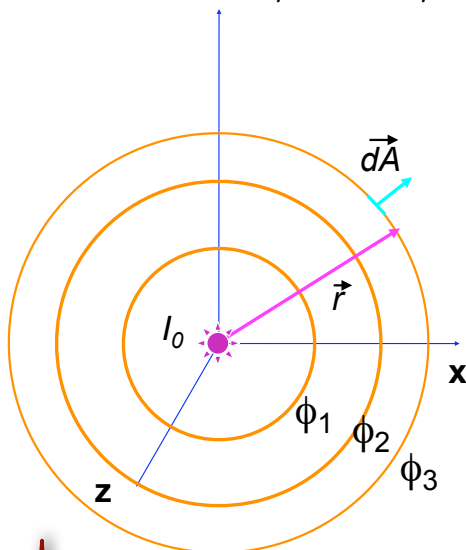


Fields

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Current Monopole Field

$$y \quad \vec{J}_r = \sigma \vec{E}_r = \sigma \frac{\partial \phi}{\partial r} \quad I_v = \int \vec{J}_r \cdot d\vec{A} = J_r 4\pi r^2$$



$$\sigma \frac{\partial \phi}{\partial r} = -\frac{I_v}{4\pi r^2} \vec{a}_r$$

$$\frac{\partial \phi}{\partial r} = -\frac{I_v}{4\pi \sigma r^2} \vec{a}_r$$

$$\phi_m = -\int \frac{I_v}{4\pi \sigma r^2} \vec{a}_r \cdot d\vec{r}$$

$$\phi_m = \frac{I_v}{4\pi \sigma r} + C$$



Fields

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Current Dipole Field

$$\phi_m = \frac{I_v}{4\pi\sigma r} + C$$

$$\phi_d = -\frac{I_v}{4\pi\sigma r_0} + \frac{I_v}{4\pi\sigma r_1}$$

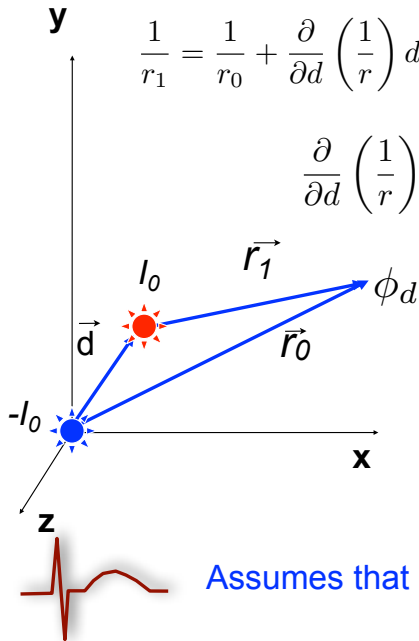
$$\frac{1}{r_1} = \frac{1}{r_0} + \frac{\partial}{\partial d} \left(\frac{1}{r} \right) d + \dots \quad \phi_d = -\frac{I_v}{4\pi\sigma r_0} + \frac{I_v}{4\pi\sigma} \left[\frac{1}{r_0} + \frac{\partial}{\partial d} \left(\frac{1}{r} \right) d \right]$$

$$\frac{\partial}{\partial d} \left(\frac{1}{r} \right) = \nabla \left(\frac{1}{r} \right) \cdot \vec{a}_d \quad \phi_d = \frac{I_0}{4\pi\sigma} \frac{\partial}{\partial d} \left(\frac{1}{r} \right) d + \dots$$

$$\phi_d = \frac{I_0}{4\pi\sigma} \left[\nabla \left(\frac{1}{r} \right) \cdot \vec{a}_d \right] d$$

$$I_0 \vec{a}_d d = \vec{p}$$

$$\phi_d = \frac{1}{4\pi\sigma} \nabla \left(\frac{1}{r} \right) \cdot \vec{p}$$



Assumes that $d \rightarrow 0$, "mathematical dipole"

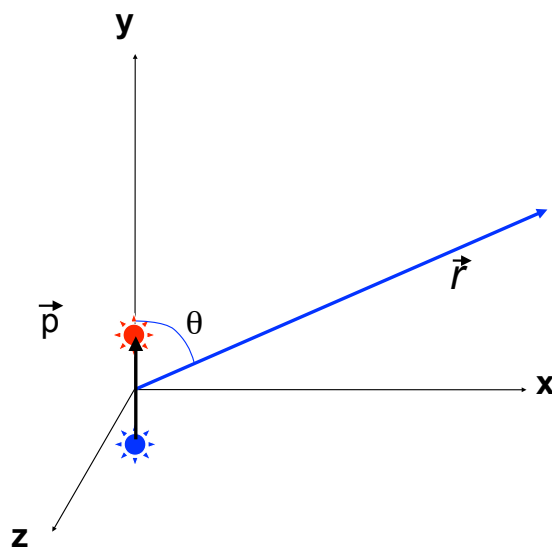
Fields

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Example

$$\phi_d = \frac{I_0}{4\pi\sigma} \nabla \left(\frac{1}{r} \right) \cdot \vec{p}$$

$$\phi_d = \frac{p \cos \theta}{4\pi\sigma r^2}$$



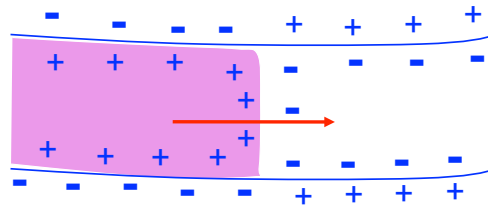
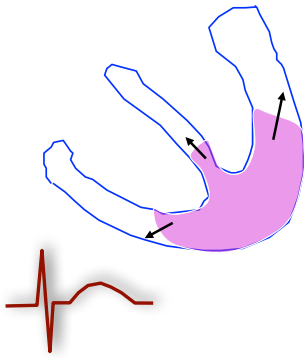
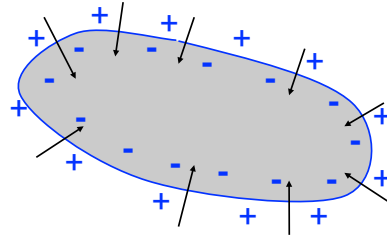
Potential from a dipole in an infinite homogeneous medium

Fields

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Why Dipoles?

- Represent bioelectric sources
 - Membrane currents
 - Coupled cells
 - Activation wavefront
 - Whole heart



Fields

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