Medical Image Analysis

CS 593 / 791

Computer Science and Electrical Engineering Dept. West Virginia University

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Discretizing the heat equation

- Mapping the image to a vector
- Boundary Conditions
- Stability





Recall : The heat equation

In 1D



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Recall : Numerical Derivatives

First order forward difference:

$$I'(x_0) \approx I(x_0 + 1) - I(x_0)$$

First order backward difference:

$$I'(x_0) \approx I(x_0) - I(x_0 - 1)$$

Second order, second centered difference:

$$I''(x_0) \approx I(x_0+1) - 2I(x_0) + I(x_0-1)$$

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Discretized 1D Heat Equation : Explicit

Using the forward difference in time we get

Explicit formulation

$$I_{x}^{t+\delta} = I_{x}^{t} + \delta(I_{x+1}^{t} - 2I_{x}^{t} + I_{x-1}^{t})$$

Explicit : Update I^t using derivatives computed at time *t*. Form a vector, **w** of image values, so that $\mathbf{w}_i = I(i)$

Discretized 1D Heat Equation : Explicit

We can rewrite the discretized heat equation as the system of linear equations:

$$w_i^{t+\delta} = [\delta, 1 - 2\delta, \delta] \begin{bmatrix} w_{i-1}^t \\ w_i^t \\ w_{i+1}^t \end{bmatrix}$$

This is equivalent to

$$w_i^{t+\delta} = [0, \delta, 1 - 2\delta, \delta, 0] \begin{bmatrix} w_{i-2}^t \\ w_{i-1}^t \\ w_i^t \\ w_{i+1}^t \\ w_{i+2}^t \end{bmatrix}$$

We can continue padding the row vector of coefficients with 0 entries until...

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Discretized 1D Heat Equation : Explicit

$$w_i^{t+\delta} = [0, \dots, 0, \delta, 1 - 2\delta, \delta, 0, \dots, 0] \mathbf{w}^t$$
$$= \mathbf{a}_i \mathbf{w}^t$$

Where $(1 - 2\delta)$ is in the i-th column, since it multiplies w_i^t . We can write the whole system of equations by forming a matrix **A** whose i-th row is **a**_i

$$\mathbf{w}^{t+1} = \mathbf{A}\mathbf{w}^t$$

A is a tridiagonal matrix.

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Discretized 1D Heat Equation : Explicit

$$w_i^{t+\delta} = [0, \dots, 0, \delta, 1 - 2\delta, \delta, 0, \dots, 0] \mathbf{w}^t$$
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$$\mathbf{w}^{t+1} = \mathbf{A}\mathbf{w}^t$$

A is a tridiagonal matrix.

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Discretized 2D Heat Equation : Explicit

Recall the 2D heat equation

$$\frac{\partial I}{\partial t} = \frac{\partial^2 I}{\partial x^2} + \frac{\partial^2 I}{\partial y^2}$$

Using the forward difference in time we get

Explicit formulation

$$I_{x,y}^{t+\delta} = I_{x,y}^{t} + \delta(I_{x+1,y}^{t} - 4I_{x,y}^{t} + I_{x-1,y}^{t} + I_{x,y+1}^{t} + I_{x,y-1}^{t})$$

Update I^t using derivatives computed at time t.

Discretized 2D Heat Equation : Implicit

Recall the heat equation

$$\frac{\partial I}{\partial t} = \frac{\partial^2 I}{\partial x^2} + \frac{\partial^2 I}{\partial y^2}$$

Using the backward difference in time we get

Explicit formulation

$$I_{x,y}^{t+\delta} = I_{x,y}^{t} + \delta(I_{x+1,y}^{t+\delta} - 4I_{x,y}^{t+\delta} + I_{x-1,y}^{t+\delta} + I_{x,y+1}^{t+\delta} + I_{x,y-1}^{t+\delta})$$

Implicit : Update I^t using derivatives computed at time $t + \delta$.

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Mapping the image to a vector

2D image indices to 1D image index

Map 2d coordinates of I(x, y)

(0,0)	(1,0)	(2,0)
(0,1)	(1,1)	(2,1)
(0,2)	(1,2)	(2,2)

to 1d coordinates of w(i)

0	3	6
1	4	7
2	5	8

The coordinate transformation is given by

$$i(x,y) = nx + y$$

for an $n \times n$ image.

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Writing central differences in 1D vector form

For the coordinate transformation function

$$i(x,y) = nx + y$$

- If $I(x, y) \to w(i)$, then $I(x, y + 1) \to w(i + 1)$, since i(x, y + 1) = nx + y + 1 = i(x, y) + 1.
- If $I(x, y) \to w(i)$, then $I(x + 1, y) \to w(i + n)$, since i(x + 1, y) = n(x + 1) + y = i(x, y) + n.

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Writing central differences in 1D vector form

For the coordinate transformation function

$$i(x,y) = nx + y$$

- If $I(x, y) \to w(i)$, then $I(x, y + 1) \to w(i + 1)$, since i(x, y + 1) = nx + y + 1 = i(x, y) + 1.
- If $l(x, y) \to w(i)$, then $l(x + 1, y) \to w(i + n)$, since i(x + 1, y) = n(x + 1) + y = i(x, y) + n.

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Writing central differences in 1D vector form

For the coordinate transformation function

$$i(x,y) = nx + y$$

- If $I(x, y) \to w(i)$, then $I(x, y + 1) \to w(i + 1)$, since i(x, y + 1) = nx + y + 1 = i(x, y) + 1.
- If *I*(*x*, *y*) → *w*(*i*), then *I*(*x* + 1, *y*) → *w*(*i* + *n*), since *i*(*x* + 1, *y*) = *n*(*x* + 1) + *y* = *i*(*x*, *y*) + *n*.

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Writing central differences in 1D vector form

For the coordinate transformation function

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• If
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• If $I(x, y) \to w(i)$, then $I(x + 1, y) \to w(i + n)$, since i(x + 1, y) = n(x + 1) + y = i(x, y) + n.

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Writing central differences in 1D vector form

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• If $I(x, y) \to w(i)$, then $I(x + 1, y) \to w(i + n)$, since i(x + 1, y) = n(x + 1) + y = i(x, y) + n.

Perona-Malik

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Writing central differences in 1D vector form

For the coordinate transformation function

$$i(x,y) = nx + y$$

So,

$$\frac{\partial l^2}{\partial y^2}(x,y) \approx l(x,y+1) - 2l(x,y) + l(x,y-1)$$
$$\approx w(i+1) - 2w(i) + w(i-1)$$

and

$$\frac{\partial l^2}{\partial x^2}(x,y) \approx l(x+1,y) - 2l(x,y) + l(x-1,y)$$
$$\approx w(i+n) - 2w(i) + w(i-n)$$

Writing central differences in 1D vector form

For the coordinate transformation function

$$i(x,y) = nx + y$$

So,

$$\frac{\partial l^2}{\partial y^2}(x,y) \approx l(x,y+1) - 2l(x,y) + l(x,y-1)$$
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and

~

$$\frac{\partial l^2}{\partial x^2}(x,y) \approx l(x+1,y) - 2l(x,y) + l(x-1,y)$$
$$\approx w(i+n) - 2w(i) + w(i-n)$$

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Writing difference equations in matrix form

The implicit formulation of the heat equation involves solving n^2 simultaneous equations:

What to do when $w(i \pm 1)$ or $w(i \pm n)$ falls outside the image

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Boundary Conditions

Constant Boundary Value



$$(x < 0)$$
 or $(x > n) \rightarrow I(x) = c$
For $c = 0$:

$$I_{xx}(0)\approx -2I(0)+I(1)$$

Boundary Conditions

Constant Boundary Slope



Fixing the slope at zero (adiabatic) gives $(x < 0) \rightarrow I(x) = I(0)$ $(x > n) \rightarrow I(x) = I(n)$

 $I_{xx}(0)\approx -I(0)+I(1)$

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Boundary Conditions

Periodic Boundary Conditions



Perona-Malik

Boundary Conditions

Reflective Boundary Conditions



$$(x < 0) \rightarrow I(x) = I(-x)$$

(x > n) $\rightarrow I(x) = I(2n - x)$

$$I_{xx}(0)\approx -2I(0)+2I(1)$$

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Stability of the explicit 1D heat equation

The 1D heat equation, $I_t = I_{xx}$, has solution $I(x, t) = e^{-t} \cos(x)$. This corresponds to the problem with initial condition $I(x, 0) = \cos(x)$.

Discretize only in time (forward)

Observe that $I_{xx}(x, t) = -e^{-t}\cos(x) = -I(x, t)$

$$\frac{l^{t+\delta}-l^t}{\delta}=-l^t$$

$$I^{t+\delta} = I^t - \delta I^t$$

Convergence criterion : ratio test

The sequence I^t is convergent if

$$\lim_{\to\infty}\left|\frac{I^{t+\delta}}{I^t}\right|<1$$

The explicit equation we formed earlier

$$I^{t+\delta} = I^t - \delta I^t$$

has convergence criterion

$$\left|\frac{I^{t+\delta}}{I^t}\right| = |1-\delta| < 1$$

This is satisfied for $0 < \delta < 2$. (Only conditionally convergent.)

Stability of the implicit 1D heat equation

Discretize only in time (backward)

$$\frac{l^{t+\delta}-l^t}{\delta} = -l^{t+\delta}$$

$$I^{t+\delta} = I^t - \delta I^{t+\delta}$$

The implicit equation has convergence criterion

$$\left|\frac{I^{t+\delta}}{I^t}\right| = \left|\frac{1}{1+\delta}\right| < 1$$

This is satisfied for $\delta > 0$.

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Stability

In general, it can be shown that

- Explicit methods are conditionally stable.
- Implicit methods are unconditionally stable.

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Outline



Discretizing the heat equation

Perona-Malik

- Introduction
- Weaknesses of the standard scale-space paradigm
- Inhomogeneous diffusion
- Properties of inhomogeneous diffusion
- Next Class

Introduction



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The need for multiscale image representations: Details in images should only exist over certain ranges of scale.





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Definition: a family of images, I(x, y, t), where

- The scale-space parameter is *t*.
- I(x, y, 0) is the original image.
- Increasing *t* corresponds to coarser resolutions.

I(x, y, t) can be generated by convolving with wider Gaussian kernels as *t* increases, or equivalently, by solving the heat equation.

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Introduction

Earlier Scale-space properties

- Causality: coarse details are "caused" by fine details.
- New details should not arise in coarse scale images.
- Smoothing should be homogeneous and isotropic.

This paper will challenge the last property, and propose a more useful scale-space definition.

The new scale-space will be shown to obey the causality property.

Weaknesses of the standard scale-space paradigm

Lost Edge Information

- Edge location is not preserved across the scale space.
- Edge crossings may disappear.
- Region boundaries are blurred.



Gaussian blurring is a local averaging operation. It does not respect natural boundaries.

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Weaknesses of the standard scale-space paradigm



- **Def:** Scale spaces generated by a linear filtering operation.
 - Nonlinear filters, such as the median filter, can be used to generate scale-space images.
 - Many nonlinear filters violate one of the scale-space conditions.

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Weaknesses of the standard scale-space paradigm

New Criteria

- Causality.
- Immediate localization : edge locations remain fixed.
- Piecewise Smoothing : permit discontinuities at boundaries.

At all scales the image will consist of smooth **regions** separated by **boundaries** (edges).

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Diffusion equation

$$\frac{\partial I}{\partial t} = \operatorname{div}(c(x, y, t) \nabla I)$$

The diffusion coefficient, c(x, y, t) controls the degree of smoothing at each point in *I*.

The basic idea:

Setting c(x, y, t) = 0 at region boundaries, and c(x, y, t) = 1 at region interior will encourage intraregion smoothing, and discourage interregion smoothing.

Diffusion equation

By the chain rule:

$$\begin{aligned} \frac{\partial I}{\partial t} &= \operatorname{div} \left(\begin{array}{c} c(x, y, t) \frac{\partial I}{\partial x} \\ c(x, y, t) \frac{\partial I}{\partial y} \end{array} \right) \\ &= \frac{\partial c}{\partial x} \frac{\partial I}{\partial x} + c(x, y, t) \frac{\partial^2 I}{\partial x^2} + \frac{\partial c}{\partial y} \frac{\partial I}{\partial y} + c(x, y, t) \frac{\partial^2 I}{\partial y^2} \\ &= c(x, y, t) \nabla^2 I + \nabla c \cdot \nabla I \end{aligned}$$

Notation

The paper uses the symbol Δ to represent the Laplacian. $\Delta I = \nabla^2 I = \operatorname{div}(\nabla I)$

Conduction coefficient

What properties would he like c(x, y, t) to have?

- c = 1 at interior of a region.
- c = 0 at boundary of a region.
- c should be nonnegative everywhere.

Since c(x, y, t) depends on edge information, we need an edge descriptor, E(x, y, t), to compute *c*.

Notation

When written as a function of the edge descriptor, the authors use the symbol g() for conduction coefficient.

Edge Estimate (or Edge Descriptor)

E(x, y, t) should convey the following information:

- Location.
- Magnitude (contrast across edge).
- Direction.

and obey the following properties:

- E(x, y, t) = 0 at region interior.
- E(x, y, t) = Ke(x, y, t) at region boundaries.

K is the contrast, $\mathbf{e}(x, y, t)$ is perpendicular to the edge. $\nabla I(x, y, t)$ has these properties, and is a useful edge estimator.

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Maximum Principle

- The maximum and minimum intensities in the scale-space image I(x, y, t) occur at t = 0 (the finest scale image).
- Since new maxima and minima correspond to new image features, the causality requirement of scale-space can satisfied if the evolution equation obeys the maximum principle.
- We will make some less rigorous observations concerning causality...

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Properties of inhomogeneous diffusion

Maximum Principle



- Solving the heat equation is equivalent to convolution.
- Convolution is a local averaging operation.
- Averaging is bounded by the values being averaged.

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Maximum Principle

For the Perona-Malik equation

$$\frac{\partial I}{\partial t} = c(x, y, t) \nabla^2 I + \nabla c \cdot \nabla I$$

Note that at local minima $\nabla I = \mathbf{0}$ and we are evolving by the original heat equation.

It can be shown that this general class of PDEs obeys the maximum principle.

We will also inspect a maximum principle for the discretized equations.

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Edge Enhancement

Inhomogeneous diffusion may actually enhance edges, for a certain choice of c(x, y, t).

1D example:

Let
$$s(x) = \frac{\partial I}{\partial x}$$
, and $\phi(s) = g(s)s = g(I_x)I_x$.

The 1D heat equation becomes

$$l_t = \frac{\partial}{\partial x} (g(I_x)I_x) = \frac{\partial}{\partial x} \phi(s)$$

by chain rule = $\frac{\partial \phi}{\partial s} \frac{\partial s}{\partial x}$
= $\phi'(s)I_{xx}$

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Edge Enhancement

We are interested in the rate of change of edge slope with respect to time.

$$\frac{\partial}{\partial t}(I_x) = \frac{\partial}{\partial x}(I_t) \text{ if I is smooth} \\ = \frac{\partial}{\partial x}(\phi'(s)I_{xx}) \\ = \phi''(s)I_{xx}^2 + \phi'(s)I_{xxx}$$

Edge Enhancement



$$\frac{\partial}{\partial t}(I_{\mathsf{x}}) = \phi''(s)I_{\mathsf{x}\mathsf{x}}^2 + \phi'(s)I_{\mathsf{x}\mathsf{x}\mathsf{x}}$$

For a step edge with $l_x > 0$ look at the inflection point, *p*.

Observe that $I_{xx}(p) = 0$, and $I_{xxx}(p) < 0$.

$$rac{\partial}{\partial t}(I_{\mathbf{x}})(\mathbf{p}) = \phi'(\mathbf{s})I_{\mathbf{x}\mathbf{x}\mathbf{x}}(\mathbf{p})$$

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The sign of this quantity depends only on $\phi'(s)$.

Edge Enhancement

At the inflection point:

$$rac{\partial}{\partial t}(I_{\mathbf{x}})(\mathbf{p}) = \phi'(\mathbf{s})I_{\mathbf{x}\mathbf{x}\mathbf{x}}(\mathbf{p})$$

- If φ'(s) > 0, then ∂/∂t(I_x)(p) < 0 (slope is decreasing).
- If $\phi'(s) < 0$, then $\frac{\partial}{\partial t}(I_x)(p) > 0$ (slope is increasing).

Since $\phi(s) = g(s)s$, selecting the function g(s) determines which edges of smoothed and which are sharpened.



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Perona and Malik suggest two possible functions

$$g(|\nabla I|) = e^{-(\frac{||\nabla I||}{\kappa})^2}$$

$$g(|
abla I|) = rac{1}{1+(rac{||
abla I||}{K})^{1+lpha}} \quad (lpha > 0)$$

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Next Class

We will continue to discuss the Perona-Malik paper, looking at parameter setting and implementation details.