Math 6630: Numerical Solutions of Partial Differential Equations Interpolation with Fourier Series See Hesthaven, S. Gottlieb, and D. Gottlieb 2007, Chapters 2-3, Canuto et al. 2011, Chapter 2.1,

Shen, Tang, and Wang 2011, Chapter 2

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#### Fourier Series approximations

We have established that Fourier series approximations  $u_N$ ,

$$u_N(x) = \sum_{|k| \leq N} \hat{u}_k \phi_k(x), \qquad \phi_k(x) = \frac{1}{\sqrt{2\pi}} e^{ikx}, \qquad \hat{u}_k = \langle u, \phi_k \rangle,$$

have orders of convergence that depend on the smoothness of u:

$$u \in H_p^s \implies ||u - u_N||_{L^2} \leq N^{-s} ||u||_{H^s}.$$

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#### $Smoothness \implies Compressibility$

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#### Smoothness $\implies$ Compressibility

One major outstanding question is *how* we actually compute  $\hat{u}_k$  in practice.

# Quadrature

The expansion coefficients require computing an integral,

$$\widehat{u}_k = \frac{1}{\sqrt{2\pi}} \int_0^{2\pi} u(x) e^{-ikx} \mathrm{d}x,$$

which can be intractable, even if an explicit formula for u is available.

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A standard recourse is to approximate the integral with quadrature:

$$\frac{1}{\sqrt{2\pi}} \int_0^{2\pi} u(x) e^{-ikx} dx \approx \sum_{j=1}^M w_{k,j} u(x_j), \quad w_{k,j} = \frac{\sqrt{2\pi}}{M} e^{-ikx_j}, \quad x_j = \frac{2\pi(j-1)}{M},$$

where we have made particular choices:

- $x_j$  are equispaced on  $[0, 2\pi]$  for  $j \in [M]$
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- We'll also assume that M = 2N + 1. (Quadrature nodes = expansion coefficients) Note that this is just the trapezoid rule on  $[0, 2\pi]$  with periodic boundary conditions.

One can make other choices, but these choices are most convenient for discussing the major concepts surrounding theory and computation.

## Quadrature as linear algebra

$$\widehat{u}_k \approx \widetilde{u}_k \coloneqq \sum_{j=1}^M w_{k,j} u(x_j), \qquad w_{k,j} = \frac{\sqrt{2\pi}}{M} e^{-ikx_j}, \qquad x_j = \frac{2\pi(j-1)}{M}.$$

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A simple implementation quadrature of amounts to matrix-vector algebra:

$$oldsymbol{u}\coloneqq egin{pmatrix} u(x_1)\ u(x_2)\ dots\ u(x_M) \end{pmatrix}, \quad \widetilde{oldsymbol{u}}\coloneqq egin{pmatrix} \widetilde{u}_{-N}\ \widetilde{u}_{-N+1}\ dots\ \widetilde{u}_{N} \end{pmatrix} \implies \widetilde{oldsymbol{u}}=\widetilde{oldsymbol{V}}^{oldsymbol{*}}oldsymbol{u},$$

where  $\widetilde{m{V}}^{*}$  is the conjugate transpose of  $\widetilde{m{V}}$ , which in turn is given by,

$$\widetilde{\boldsymbol{V}} = \sqrt{\frac{2\pi}{M}} \boldsymbol{V}, \qquad \boldsymbol{V} = \begin{pmatrix} | & | & | & | \\ \boldsymbol{v}_{-N} & \boldsymbol{v}_{-N+1} & \cdots & \boldsymbol{v}_{N} \\ | & | & | \end{pmatrix}, \qquad \boldsymbol{v}_{k} = \sqrt{\frac{2\pi}{M}} \phi_{k}(\boldsymbol{x}),$$

and  $\boldsymbol{x} = (x_1, x_2, \dots, x_M)^T$ .

# Duality of Fourier Series with quadrature

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A somewhat straightforward computation shows:

$$\langle \boldsymbol{v}_{\ell}, \boldsymbol{v}_{k} \rangle = \frac{1}{M} \sum_{j=1}^{M} e^{i(\ell-k)2\pi(j-1)/M} = \frac{1}{M} \sum_{j=0}^{M-1} \left( e^{i(\ell-k)2\pi/M} \right)^{j},$$

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Thus, in particular if  $\ell = k$  then  $\langle \boldsymbol{v}_{\ell}, \boldsymbol{v}_{k} \rangle = 1$ , and for  $\ell \neq k$  and  $|\ell - k| \leq M - 1$ :

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I.e.,  $\{v_{k}\}_{|k| \leq N}$  are orthonormal vectors.

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I.e.,  $\{v\}_{|k| \leqslant N}$  are orthonormal vectors.

This shows the important property that V is a unitary matrix:

$$V^*V = I \implies V^{-1} = V^*.$$

# Inverting the Fourier Series

Putting everything together:

$$\widetilde{\boldsymbol{u}} = \widetilde{\boldsymbol{V}}^* \boldsymbol{u}, \qquad \qquad \widetilde{\boldsymbol{V}} = \sqrt{\frac{2\pi}{M}} \boldsymbol{V}, \qquad \qquad \boldsymbol{V}^{-1} = \boldsymbol{V}^*$$

This implies that:

$$\boldsymbol{u} = \left(\widetilde{\boldsymbol{V}}^*\right)^{-1} \widetilde{\boldsymbol{u}} = \sqrt{\frac{M}{2\pi}} \left(\boldsymbol{V}^*\right)^{-1} \widetilde{\boldsymbol{u}} = \sqrt{\frac{M}{2\pi}} \boldsymbol{V} \widetilde{\boldsymbol{u}} = \frac{M}{2\pi} \widetilde{\boldsymbol{V}} \widetilde{\boldsymbol{u}}.$$

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I.e., the map between  $oldsymbol{u}$  and  $\widetilde{oldsymbol{u}}$  is invertible and quite explicit:

$$\widetilde{\boldsymbol{u}} = \widetilde{\boldsymbol{V}}^* \boldsymbol{u}, \qquad \qquad \boldsymbol{u} = \frac{M}{2\pi} \widetilde{\boldsymbol{V}} \widetilde{\boldsymbol{u}}.$$

This invertible map is called the Discrete Fourier Transform (DFT). As a consequence of V being unitary, we have also shown that the DFT is a (scaled) isometry,

$$\int_{0}^{2\pi} |u(x)|^{2} \mathrm{d}x \approx \frac{2\pi}{M} \|\boldsymbol{u}\|_{2}^{2} = \|\widetilde{\boldsymbol{u}}\|_{2}^{2},$$

which is the discrete analogue of Parseval's identity.

#### The Fast Fourier Transform, I

The inverse/DFT is relatively expensive:

$$\boldsymbol{u} \xrightarrow{\mathcal{O}(M^2)} \widetilde{\boldsymbol{V}}^* \boldsymbol{u}, \qquad \qquad \widetilde{\boldsymbol{u}} \xrightarrow{\mathcal{O}(M^2)} \frac{M}{2\pi} \widetilde{\boldsymbol{V}}.$$

One of the most well-known algorithms is the *fast Fourier transform*, which is a fast algorithm for accomplishing the particular matrix-vector multiplication  $\tilde{V}^* u$ .

It is simpler to explain the basic idea if M is even, in which case we have:

$$\frac{M}{\sqrt{2\pi}}\tilde{u}_k = \sum_{j=1}^M u(x_j)e^{-ikx_j} = \sum_{j=1}^M u(x_j)e^{-ik2\pi(j-1)/M}$$
$$= \sum_{j=1}^{M/2} u(x_{2j})e^{-ik2\pi 2(j-1)/M} + \sum_{j=1}^{M/2} u(x_{2j-1})e^{-ik2\pi(2j-1)/M}$$

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$$\begin{split} \frac{M}{\sqrt{2\pi}} \widetilde{u}_k &= \sum_{j=1}^M u(x_j) e^{-ikx_j} = \sum_{j=1}^M u(x_j) e^{-ik2\pi(j-1)/M} \\ &= \sum_{j=1}^{M/2} u(x_{2j}) e^{-ik2\pi 2(j-1)/M} + \sum_{j=1}^{M/2} u(x_{2j-1}) e^{-ik2\pi(2j-1)/M} \\ &= \sum_{j=1}^{M/2} u(x_{2j}) e^{-ik2\pi 2(j-1)/M} + e^{ik2\pi/M} \sum_{j=1}^{M/2} u(x_{2j-1}) e^{-ik2\pi(2j-2)/M}. \end{split}$$

Note that the last two sums are M/2-point DFT coefficients associated with half the data (either at  $x_{2j}$  or at  $x_{2j-1}$ ).

I.e., with some book-keeping, we can compute the  $M\mbox{-}{\rm point}$  DFT using 2  $M/2\mbox{-}{\rm point}$  DFT's.

# The Fast Fourier Transform, II

This logic can be repeated, showing that actually we can compute the M-point DFT using J(M/J)-point DFT's, where J is a power of two. This yields the simplest, radix 2 fast Fourier transform (FFT) algorithm.



Through this divide-and-conquer strategy, an M-point DFT that naively requires  $\mathcal{O}(M^2)$  complexity can be accomplished in  $\mathcal{O}(M \log M)$  time.

## Interpolation

$$\widetilde{\boldsymbol{u}} = \widetilde{\boldsymbol{V}}^* \boldsymbol{u}, \qquad \qquad \boldsymbol{u} = \frac{M}{2\pi} \widetilde{\boldsymbol{V}} \widetilde{\boldsymbol{u}}.$$

We have introduced the DFT via quadrature, but an alternative and illustrative viewpoint is *interpolation*.

Note that the coefficients  $\widetilde{u}$  are determined by the conditions,

$$\frac{M}{2\pi}\widetilde{\boldsymbol{V}}\widetilde{\boldsymbol{u}} = \boldsymbol{u} \implies \begin{pmatrix} | & | & | & | \\ \phi_{-N} & \phi_{-N+1} & \cdots & \phi_{N} \\ | & | & | \end{pmatrix} \widetilde{\boldsymbol{u}} = \boldsymbol{u}.$$

Note that these are "just" interpolation conditions for the  $\tilde{u}$  at the data points  $x_j$ ,  $j \in [M]$ .

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Note that these are "just" interpolation conditions for the  $\tilde{u}$  at the data points  $x_j$ ,  $j \in [M]$ .

Hence,  $u_N(x) = \sum_{|k| \leq N} \widetilde{u}_k \phi_k(x)$  interpolates the data  $\boldsymbol{u}$ .

We have already concluded that this interpolation problem is *unisolvent*. Hence, there are cardinal basis functions  $\ell_j(x)$ ,  $j \in [M]$  such that,

$$u_N(x) = \sum_{j \in [M]} u(x_j)\ell_j(x), \qquad \qquad \ell_j(x_r) = \delta_{j,r}.$$

I.e.,  $u_N$  has a Lagrange form.

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# Cardinal Lagrange basis

The cardinal Lagrange functions yield insight into the interpolation process.



Figure 2.3 of Hesthaven, S. Gottlieb, and D. Gottlieb 2007

Note that interpolation implies

$$u(x) \in V_N \coloneqq \operatorname{span}\left\{e^{ikx}\right\}_{|k| \leqslant N} \implies I_N u \coloneqq \sum_{|k| \leqslant N} \widetilde{u}_k \phi_k(x) = u(x).$$

# Aliasing

The fact that our DFT is an interpolation process reveals a significant issue that we must be cognizant of: aliasing error.



Figure 2.7 of Hesthaven, S. Gottlieb, and D. Gottlieb 2007

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So, for example, even if  $\langle e^{i\ell x}, \phi_k(x) \rangle = 0$  for  $\ell > N$ , it's possible that  $I_N u \neq 0$ .

I.e., the interpolation/DFT procedure *is* a projection operator, it's just an oblique one.

# Aliasing error

Aliasing is not just an academic curiosity: with  $P_N$  the  $L^2$ -orthogonal projection operator onto

$$V_N = \operatorname{span}\left\{e^{ikx}\right\}_{|k| \leqslant N},$$

recall that  $u \in H_p^s$  implies that  $||u - P_N u||_{L^2} \leq N^{-s}$ .

Ok, but what about  $I_N u$ ?

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#### Ok, but what about $I_N u$ ?

The main strategy to understanding this is to estimate the aliasing error. Note that for the  $L^2$  norm,

$$\|u - I_N u\| = \| (u - P_N u) + (P_N u - I_N u) \| \leq \|u - P_N u\| + \|P_N u - I_N u\|$$
  
=:  $\|u - P_N u\| + \|A_N u\|$ ,

where we have defined the aliasing error  $A_N u$ .

$$A_N u = P_N u - I_N u$$

The following observations are crucial:

- If  $u \in V_N$ , then  $I_N u = P_N u = u$ , so  $A_N u = 0$ . Therefore,  $A_N u = A_N (I - P_N) u$ . The aliasing error is only affected by truncation error.

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- We know that  $(I P_N)u$  is small. Therefore, if  $A_N$  is "behaves well", then  $A_N u$  will be small.

The truncation error is small, but does  $A_N$  amplify small inputs?

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The truncation error is small, but does  $A_N$  amplify small inputs?

-  $A_N$  is well-behaved: for  $|k| \leq N$ ,  $I_N A_N e^{i(k+(2N+1))x} = e^{ikx}, \qquad P_N e^{i(k+(2N+1))x} = 0$ 

and thus in particular,

$$u = \sum_{\substack{|k| \leq N}} \hat{u}_k \phi_k(x) \implies \tilde{u}_k = \sum_{\ell \in \mathbb{Z}} \hat{u}_{k+\ell(2N+1)}.$$

 $A_N$  does not amplify small inputs.

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The truncation error is small, but does  $A_N$  amplify small inputs?  $I_{N} \hat{\mathcal{U}}_{V} q_{t} - \hat{\mathcal{U}}_{V} q_{\ell}$   $I_{N} \hat{\mathcal{Q}}_{\ell} - e^{i \ell_{X}}$   $|\ell| \leq N$ 

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 $A_N$  does not amplify small inputs.  $|\ell| \leq N \leq \langle \mathbf{I}_N \mathbf{V}, q_\ell \rangle = \sum_{\mathbf{V} \in \mathcal{I}} \hat{\mathbf{V}}_{\mathbf{V}} \langle \mathbf{I}_N \phi_{\mathbf{V}}, q_\ell \rangle$ Therefore, if  $\hat{u}_{k+\ell(2N+1)}$  decays quickly for large  $|\ell|$ , then we can expect the aliased coefficients  $\widetilde{u}_k$  to be "close" to  $\widehat{u}_k$ . = Z úk & K, X+ (2N+1)g While we have only discussed the high-level ideas, going through the details produces the following estimate:

# Theorem Assume $u \in H_p^s$ with s > 1/2. Then $\|u - I_N u\|_{L^2} \lesssim N^{-s} \|u\|_{H^s}$ . $\|u - I_N u\|_{H^r} \lesssim N^{-(s-r)} \|u\|_{H^s}, \qquad r < s.$

Note that this is exactly the asymptotic behavior for the exact orthogonal projector  $P_N$ . Thus, one can expect the DFT to produce good results.

# The DFT in practice



Figure 2.4 of Hesthaven, S. Gottlieb, and D. Gottlieb 2007

$$u(x) = \frac{3}{5 - 4\cos x}$$

# The DFT in practice



Figure 2.4 of Hesthaven, S. Gottlieb, and D. Gottlieb 2007

 $u(x) = \sin(x/2)$ 



# References I

Canuto, Claudio et al. (2011). *Spectral Methods: Fundamentals in Single Domains*. 1st ed. 2006. Corr. 4th printing 2010 edition. Berlin ; New York: Springer. ISBN: 978-3-540-30725-9.

Hesthaven, Jan S., Sigal Gottlieb, and David Gottlieb (2007). *Spectral Methods for Time-Dependent Problems*. Cambridge University Press. ISBN: 0-521-79211-8.

Shen, Jie, Tao Tang, and Li-Lian Wang (2011). *Spectral Methods: Algorithms, Analysis and Applications*. Springer Science & Business Media. ISBN: 978-3-540-71041-7.