# Math 6630: Numerical Solutions of Partial Differential Equations Approximation with Fourier Series 

See Hesthaven, P. S. Gottlieb, and D. Gottlieb 2007, Chapters 1-2,
Canuto et al. 2011, Chapters 2.1, 5.1,
Shen, Tang, and Wang 2011, Chapter 2

Akil Narayan ${ }^{1}$
${ }^{1}$ Department of Mathematics, and Scientific Computing and Imaging (SCI) Institute University of Utah

March 13, 2023

High order approximations
We are now very familiar with our rather standard approximation to $u_{x x}$ on an equidistant grid:

$$
D_{0} u_{j}^{n} \approx u_{x x}+\mathcal{O}\left(h^{2}\right)
$$

Note that the $h^{2}$ truncation error is a direct result of our choice of 3-point stencil.
Using more points in the stencil allows us to attain higher order truncation errors.

$$
\frac{1}{12 h^{2}}\left[-u_{j-2}^{n}+16 u_{j-1}^{n}-30 u_{j}^{n}+16 u_{j+1}^{n}-u_{j+2}^{n}\right] \approx u_{x x}+\mathscr{g}^{h^{2}} \cdot o\left(h^{4}\right)
$$

## High order approximations

We are now very familiar with our rather standard approximation to $u_{x x}$ on an equidistant grid:

$$
D_{0} u_{j}^{n} \approx u_{x x}+\mathcal{O}\left(h^{2}\right) \quad(k=1)
$$

Note that the $h^{2}$ truncation error is a direct result of our choice of 3-point stencil.
Using more points in the stencil allows us to attain higher order truncation errors.

$$
\frac{1}{12 h^{2}}\left[-u_{j-2}^{n}+16 u_{j-1}^{n}-30 u_{j}^{n}+16 u_{j+1}^{n}-u_{j+2}^{n}\right] \approx u_{x x}+\varnothing^{\kappa^{K 2}} . \quad(1 /=2)
$$

In general, using $2 k+1$ points allows us to achieve $\mathcal{O}\left(h^{2 k}\right)$ LTE.
Why stop here? Why not take $k$ as large as possible?
This requires a stencil spreading over the whole domain, globally coupling all degrees of freedom.

Is it worth it?

## Fourier Series, I

Before solving differential equations, let's answer some basic approximation theory questions first.

The simplest example of an approximation scheme that globally couples all degrees of freedom is a Fourier Series.

## Fourier Series, I

Before solving differential equations, let's answer some basic approximation theory questions first.

The simplest example of an approximation scheme that globally couples all degrees of freedom is a Fourier Series.

Consider a given $u:[0,2 \pi] \rightarrow \mathbb{C}$, which we represent as a sum of complex exponentials,

$$
u(x) \approx \sum_{k \in \mathbb{Z}} \widehat{u}_{k} \phi_{k}(x), \quad \quad \phi_{k}(x)=\frac{1}{\sqrt{2 \pi}} e^{i k x}
$$

The most straightforward strategy to identify $\widehat{u}_{k}$ is to choose them to minimize a loss,

$$
\widehat{u}_{k}=\underset{\widehat{u}_{k}, k \in \mathbb{Z}}{\arg \min }\left\|u(x)-\sum_{k \in \mathbb{Z}} \widehat{u}_{k} \phi_{k}(x)\right\|_{2}^{2},
$$

where we have introduced the norm and a corresponding inner product,

$$
\langle f, g\rangle:=\int_{0}^{2 \pi} f(x) \bar{g}(x) \mathrm{d} x, \quad\|f\|_{2}^{2}:=\langle f, f\rangle
$$

where $\bar{z}$ is the complex conjugate of $z .{ }^{1}$

[^0][^1]
## Fourier Series, II

We have conveniently chosen the basis $\phi_{k}$ so that,

$$
\left\langle\phi_{k}, \phi_{\ell}\right\rangle= \begin{cases}1, & k=\ell \\ 0, & k \neq \ell\end{cases}
$$

Such basis functions are orthonormal.
There is a unique solution for the $\widehat{u}_{k}$ that minimizes the loss, and using basis orthonormality the solution has a fairly simple expression,

$$
\widehat{u}_{k}=\left\langle u, \phi_{k}\right\rangle=\frac{1}{\sqrt{2 \pi}} \int_{0}^{2 \pi} u(x) e^{-i k x} \mathrm{~d} x .
$$

## Fourier Series, II

We have conveniently chosen the basis $\phi_{k}$ so that,

$$
\left\langle\phi_{k}, \phi_{\ell}\right\rangle= \begin{cases}1, & k=\ell \\ 0, & k \neq \ell\end{cases}
$$

Such basis functions are orthonormal.
There is a unique solution for the $\widehat{u}_{k}$ that minimizes the loss, and using basis orthonormality the solution has a fairly simple expression,

$$
\widehat{u}_{k}=\left\langle u, \phi_{k}\right\rangle=\frac{1}{\sqrt{2 \pi}} \int_{0}^{2 \pi} u(x) e^{-i k x} \mathrm{~d} x .
$$

This gives us a first taste of some functional analysis: Define,

$$
L^{2}=L^{2}([0,2 \pi] ; \mathbb{C})=\left\{f:[0,2 \pi] \rightarrow \mathbb{C} \mid\|f\|_{2}^{2}<\infty\right\}
$$

Then Fourier Series representations are complete in $L^{2}$ :

$$
u \in L^{2} \Longrightarrow \lim _{N \rightarrow \infty}\left\|u(x)-\sum_{k=-N}^{N} \widehat{u}_{k} \phi_{k}(x)\right\|_{2}=0
$$

and orthonormality of the basis results in Parseval's identity,

$$
u \in L^{2} \quad \Longrightarrow \quad\|u\|_{2}^{2}=\sum_{k \in \mathbb{Z}}\left|\widehat{u}_{k}\right|^{2} .
$$

## Fourier approximation

$$
u(x) \stackrel{L^{2}}{=} \sum_{k \in \mathbb{Z}} \widehat{u}_{k} \phi_{k}(x)
$$

$$
\widehat{u}_{k}=\left\langle u, \phi_{k}\right\rangle
$$

This is all well and good, but how does this serve us computationally?
With finite storage, we have to truncate the infinite series,

$$
u(x) \approx u_{N}(x):=\sum_{|k| \leqslant N} \widehat{u}_{k} \phi_{k}(x)
$$

How well does $u_{N}$ approximate $u$ ?

## Fourier approximation

$$
u(x) \stackrel{L^{2}}{=} \sum_{k \in \mathbb{Z}} \widehat{u}_{k} \phi_{k}(x)
$$

$$
\widehat{u}_{k}=\left\langle u, \phi_{k}\right\rangle
$$

This is all well and good, but how does this serve us computationally?
With finite storage, we have to truncate the infinite series,

$$
u(x) \approx u_{N}(x):=\sum_{|k| \leqslant N} \widehat{u}_{k} \phi_{k}(x)
$$

How well does $u_{N}$ approximate $u$ ?

There is, of course, the other pesky issue that in practice we cannot actually compute $\widehat{u}_{k}$ exactly, and so must resort to additional approximations.
But let's focus on one sin at a time....

## Fourier approximation

$$
u(x) \stackrel{L^{2}}{=} \sum_{k \in \mathbb{Z}} \widehat{u}_{k} \phi_{k}(x), \quad \widehat{u}_{k}=\left\langle u, \phi_{k}\right\rangle
$$

This is all well and good, but how does this serve us computationally?
With finite storage, we have to truncate the infinite series,

$$
u(x) \approx u_{N}(x):=\sum_{|k| \leqslant N} \widehat{u}_{k} \phi_{k}(x)
$$

How well does $u_{N}$ approximate $u$ ?

There is, of course, the other pesky issue that in practice we cannot actually compute $\widehat{u}_{k}$ exactly, and so must resort to additional approximations.
But let's focus on one sin at a time....

So our question regards how compressible the infinite series is with respect to the truncation $N$ :

$$
\left\|u-u_{N}\right\|_{2}^{2} \stackrel{?}{\lesssim} h(N),
$$

for some function $h(N)$.

- $h$ decays quickly with $N \rightarrow u$ is very compressible
- $h$ decays slowly with $N \rightarrow u$ is not very compressible


## Projections

Before investigating Fourier approximation results, it's worthwhile to introduce additional concepts: Projections.

Given an operator $P: L^{2} \rightarrow V$, where $V \subset L^{2}$ is some subspace of $L^{2}$, then $P$ is a projection operator if

$$
P^{2}=P
$$

The action $u \mapsto P u$ projects $u$ onto $V$.
The action $u \mapsto(I-P) u$ projects $u$ onto some subspace $W$ such that $V \oplus W=L^{2}$.

## Projections

Before investigating Fourier approximation results, it's worthwhile to introduce additional concepts: Projections.

Given an operator $P: L^{2} \rightarrow V$, where $V \subset L^{2}$ is some subspace of $L^{2}$, then $P$ is a projection operator if

$$
P^{2}=P
$$

The action $u \mapsto P u$ projects $u$ onto $V$.
The action $u \mapsto(I-P) u$ projects $u$ onto some subspace $W$ such that $V \oplus W=L^{2}$.
For any projection operator $P$ and any $u \in L^{2}$, we have,

$$
(I-P) P u=0 .
$$

## Projections

Before investigating Fourier approximation results, it's worthwhile to introduce additional concepts: Projections.

Given an operator $P: L^{2} \rightarrow V$, where $V \subset L^{2}$ is some subspace of $L^{2}$, then $P$ is a projection operator if

$$
P^{2}=P .
$$

The action $u \mapsto P u$ projects $u$ onto $V$.
The action $u \mapsto(I-P) u$ projects $u$ onto some subspace $W$ such that $V \oplus W=L^{2}$.
For any projection operator $P$ and any $u \in L^{2}$, we have,

$$
\text { requires } \overbrace{\infty}^{\text {|tilbertiar Setting }}(I-P) P u=0 \text {. }
$$

A projection operator $P$ is orthogonal if $W \perp V$, equivalently if for every $u, v \in L^{2}$ :

$$
\begin{aligned}
u=P u+([-\rho) u & P=P^{*},
\end{aligned} \quad\left\langle P^{*} u, v\right\rangle:=\langle u, P v\rangle .
$$

## Truncation and projection

We are considering the truncation,

$$
\sum_{k \in \mathbb{Z}} \widehat{u}_{k} \phi_{k}(x) \stackrel{L^{2}}{=} u \approx u_{N}=\sum_{|k| \leqslant N} \widehat{u}_{k} \phi_{k}(x)
$$

This truncation is an orthogonal projector.

## Theorem

Define $P_{N}$ as the operator,,$N \in \mathbb{N}_{0}$

$$
P_{N} u=u_{N}=\sum_{|k| \leqslant \mathbb{Z}} \widehat{u}_{k} \phi_{k}(x), \quad u \stackrel{L^{2}}{=} \sum_{k \in \mathbb{Z}} \widehat{u}_{k} \phi_{k} .
$$

Then $P_{N}$ is an orthogonal projection operator.

$$
\left\|u-P_{N} u\right\| ?
$$

## A basic approximation estimate, I

Can we bound $\left\|u-P_{N}\right\|_{2}$ ? First note that,

$$
\left\|u-P_{N} u\right\|_{2}^{2}=\sum_{|k|>N}\left|\widehat{u}_{k}\right|^{2}
$$

## A basic approximation estimate, I

Can we bound $\left\|u-P_{N}\right\|_{2}$ ? First note that,

$$
\left\|u-P_{N} u\right\|_{2}^{2}=\sum_{|k|>N}\left|\widehat{u}_{k}\right|^{2} .
$$

Integration by parts is our friend, and note that,

$$
\begin{aligned}
\left\langle U_{1} Q_{\mid c}\right\rangle=\widehat{u}_{k} & =\frac{1}{\sqrt{2 \pi}} \int_{0}^{2 \pi} u(x) e^{-i k x} \mathrm{~d} x \\
& =\left.\frac{i}{k \sqrt{2 \pi}} u(x) e^{-i k x}\right|_{0} ^{2 \pi}-\frac{i}{k \sqrt{2 \pi}} \int_{0}^{2 \pi} u^{\prime}(x) e^{-i k x} \mathrm{~d} x
\end{aligned}
$$

## A basic approximation estimate, I

Can we bound $\left\|u-P_{N}\right\|_{2}$ ? First note that,

$$
\left\|u-P_{N} u\right\|_{2}^{2}=\sum_{|k|>N}\left|\widehat{u}_{k}\right|^{2} .
$$

Integration by parts is our friend, and note that,

$$
\begin{aligned}
\widehat{u}_{k} & =\frac{1}{\sqrt{2 \pi}} \int_{0}^{2 \pi} u(x) e^{-i k x} \mathrm{~d} x \\
& =\left.\frac{i}{k \sqrt{2 \pi}} u(x) e^{-i k x}\right|_{0} ^{2 \pi}-\frac{i}{k \sqrt{2 \pi}} \int_{0}^{2 \pi} u^{\prime}(x) e^{-i k x} \mathrm{~d} x .
\end{aligned}
$$

Note that, conveniently, the first term vanishes if $u(0)=u(2 \pi)$.
This is, of course, quite reasonable since we are approximating with periodic functions.
Note also that the remaining integral is the Fourier series coefficient for the derivative, $u^{\prime}(x)$ :

$$
u^{\prime}(x)=\sum_{|k| \in \mathbb{Z}} \widehat{u}^{\prime}{ }_{k} \phi_{k}(x), \quad{\widehat{u^{\prime}}}_{k}=\left\langle u^{\prime}, \phi_{k}\right\rangle .
$$

## A basic approximation estimate, I

Can we bound $\left\|u-P_{N}\right\|_{2}$ ? First note that,

$$
\left\|u-P_{N} u\right\|_{2}^{2}=\sum_{|k|>N}\left|\widehat{u}_{k}\right|^{2} .
$$

Integration by parts is our friend, and note that,

$$
\begin{aligned}
\widehat{u}_{k} & =\frac{1}{\sqrt{2 \pi}} \int_{0}^{2 \pi} u(x) e^{-i k x} \mathrm{~d} x \\
& =\left.\frac{i}{k \sqrt{2 \pi}} u(x) e^{-i k x}\right|_{0} ^{2 \pi}-\frac{i}{k \sqrt{2 \pi}} \int_{0}^{2 \pi} u^{\prime}(x) e^{-i k x} \mathrm{~d} x .
\end{aligned}
$$

Note that, conveniently, the first term vanishes if $u(0)=u(2 \pi)$.
This is, of course, quite reasonable since we are approximating with periodic functions.
Note also that the remaining integral is the Fourier series coefficient for the derivative, $u^{\prime}(x)$ :

$$
u^{\prime}(x)=\sum_{|k| \in \mathbb{Z}}{\widehat{u^{\prime}}}_{k} \phi_{k}(x), \quad \widehat{u}_{k}^{\prime}=\left\langle u^{\prime}, \phi_{k}\right\rangle
$$

Thus, if $u$ is periodic and $u^{\prime} \in L^{2}$ (so that $\widehat{u}_{k}$ is well-defined), then

$$
\widehat{u}_{k}=-\frac{i}{k}{\widehat{u^{\prime}}}_{k}
$$

## A basic approximation estimate, II

$$
\begin{aligned}
\left\|u-P_{N} u\right\|_{2}^{2} & =\sum_{|k|>N}\left|\widehat{u}_{k}\right|^{2} \\
\widehat{u}_{k} & =-\frac{i}{k} \widehat{u^{\prime}}{ }_{k}
\end{aligned}
$$

This very basic estimate for Fourier series coefficients implies:

$$
\begin{aligned}
\left\|u-P_{N} u\right\|_{2}^{2}=\sum_{|k|>N} \frac{1}{|k|^{2}}\left|{\widehat{u^{\prime}}}_{k}\right|^{2} \leqslant \frac{1}{N^{2}} \sum_{|k|>N}\left|{\hat{u^{\prime}}}_{k}\right|^{2} & \leqslant \frac{1}{N^{2}} \sum_{k \in \mathbb{Z}}\left|{\widehat{u^{\prime}}}_{k}\right|^{2} \\
& =\frac{1}{N^{2}}\left\|u^{\prime}\right\|_{2}^{2}
\end{aligned}
$$

where the last relation is Parseval's identity.

## A basic approximation estimate, II

$$
\begin{aligned}
\left\|u-P_{N} u\right\|_{2}^{2} & =\sum_{|k|>N}\left|\widehat{u}_{k}\right|^{2} \\
\widehat{u}_{k} & =-\frac{i}{k} \widehat{u^{\prime}}{ }_{k}
\end{aligned}
$$

This very basic estimate for Fourier series coefficients implies:

$$
\begin{aligned}
\left\|u-P_{N} u\right\|_{2}^{2}=\sum_{|k|>N} \frac{1}{|k|^{2}}\left|{\widehat{u^{\prime}}}_{k}\right|^{2} \leqslant \frac{1}{N^{2}} \sum_{|k|>N}\left|{\hat{u^{\prime}}}_{k}\right|^{2} & \leqslant \frac{1}{N^{2}} \sum_{k \in \mathbb{Z}}\left|{\widehat{u^{\prime}}}_{k}\right|^{2} \\
& =\frac{1}{N^{2}}\left\|u^{\prime}\right\|_{2}^{2}
\end{aligned}
$$

where the last relation is Parseval's identity.
We have just proven the following:

## Theorem

Suppose $u, u^{\prime} \in L^{2}$, and that $u(0)=u(2 \pi)$. Then,

$$
\left\|u-P_{N} u\right\|_{2} \leqslant \frac{1}{N}\left\|u^{\prime}\right\|_{L^{2}}
$$

## Sobolev spaces

To generalize this result, some additional notation will be helpful.

## Definition (Sobolev spaces)

Given $s \in \mathbb{N}_{0}=\{0,1, \ldots$,$\} , the ( L^{2}$ periodic) Sobolev space of functions is given by,

$$
\begin{aligned}
H_{p}^{s}([0,2 \pi] ; \mathbb{C}):=\{f:[0,2 \pi] \rightarrow \mathbb{C} \mid & f^{(k)} \in L^{2}([0,2 \pi] ; \mathbb{C}) \text { for all } 0 \leqslant k \leqslant s, \\
& \left.f^{(k)}(0)=f^{(k)}(2 \pi) \text { for all } 0 \leqslant k \leqslant s-1\right\}
\end{aligned}
$$

The norm on $H^{s}$ is defined as,

$$
\|u\|_{H^{s}}^{2}:=\sum_{k=0}^{s}\left\|u^{(k)}\right\|_{2}^{2} .
$$

Some specializations of interest:
$-s=0 \Longrightarrow H^{0}=L^{2}$
$-s>0 \Longrightarrow$ continuous functions $\not \subset H^{s}$

$$
\begin{aligned}
& f(x)=\widehat{~} \\
& f \in H^{\prime} \\
& f \notin H^{2}
\end{aligned}
$$

## Sobolev spaces

To generalize this result, some additional notation will be helpful.

## Definition (Sobolev spaces)

Given $s \in \mathbb{N}_{0}=\{0,1, \ldots$,$\} , the ( L^{2}$ periodic) Sobolev space of functions is given by,

$$
\begin{array}{ll}
H_{p}^{s}([0,2 \pi] ; \mathbb{C}):=\{f:[0,2 \pi] \rightarrow \mathbb{C} \mid & f^{(k)} \in L^{2}([0,2 \pi] ; \mathbb{C}) \text { for all } 0 \leqslant k \leqslant s, \\
\notin & \left.f^{(k)}(0)=f^{(k)}(2 \pi) \text { for all } 0 \leqslant k \leqslant s-1\right\}
\end{array}
$$

The norm on $H^{s}$ is defined as,

$$
\|u\|_{H^{s}}^{2}:=\sum_{k=0}^{s}\left\|u^{(k)}\right\|_{2}^{2} .
$$

Some specializations of interest:
$-s=0 \Longrightarrow H^{0}=L^{2}$
$-s>0 \Longrightarrow$ continuous functions $\subset H^{s}$
The parameter $s$ encodes the "amount" of smoothness that functions have, and the following inclusions hold:

$$
H_{p}^{r} \subset H_{p}^{s}, \quad r>s \geqslant 0
$$

## General approximation results

The language of Sobolev spaces is the standard language in which to technically describe convergence rates of Fourier Series approximations.

## Theorem

If $u \in H_{p}^{s}$, then

$$
\left\|u-P_{N} u\right\|_{L^{2}} \leqslant N^{-s}\|u\|_{H^{s}}
$$

Note that $s=1$ is our previous result.
In terms of degrees of freedom, $M,\left\|u-P_{N} u\right\|_{L^{2}} \lesssim(M / 2)^{-s}\|u\|_{H^{s}}$, which is fantastic for large $s$.

## General approximation results

The language of Sobolev spaces is the standard language in which to technically describe convergence rates of Fourier Series approximations.

## Theorem

If $u \in H_{p}^{s}$, then

$$
\left\|u-P_{N} u\right\|_{L^{2}} \leqslant N^{-s}\|u\|_{H^{s}}
$$

Note that $s=1$ is our previous result.
In terms of degrees of freedom, $M,\left\|u-P_{N} u\right\|_{L^{2}} \lesssim(M / 2)^{-s}\|u\|_{H^{s}}$, which is fantastic for large $s$.

Actually, something even stronger is true about Fourier approximation:

## Theorem

If $u \in H_{p}^{s}$, then for every $0 \leqslant r<s$,

$$
\left\|u-P_{N} u\right\|_{H^{r}} \leqslant N^{-(s-r)}\|u\|_{H^{s}} .
$$

This result demonstrates tradeoff between smoothness of the function versus the strength of the norm under which convergence is sought.

## References I



Canuto, Claudio et al. (2011). Spectral Methods: Fundamentals in Single Domains. 1st ed. 2006. Corr. 4th printing 2010 edition. Berlin ; New York: Springer. ISBN: 978-3-540-30725-9.


Hesthaven, Jan S., Professor Sigal Gottlieb, and David Gottlieb (2007). Spectral Methods for Time-Dependent Problems. Cambridge University Press. ISBN: 0-521-79211-8.

Shen, Jie, Tao Tang, and Li-Lian Wang (2011). Spectral Methods: Algorithms, Analysis and Applications. Springer Science \& Business Media. ISBN: 978-3-540-71041-7.


[^0]:    ${ }^{1}$ We are mostly interested in real-valued functions, so the introduction of complex arithmetic is somewhat artificial here. We could write the basis as real-valued $\sin k x$ and $\cos k x$ functions with real coefficients. This achieves the same results but uses somewhat more technical formulas.

[^1]:    A. Narayan (U. Utah - Math/SCI)

    Math 6630: Approximation with Fourier Series

