



Multiscale models of integrins and cellular adhesions

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Abstract

Computational models of integrin-based adhesion complexes have revealed important insights into the mechanisms by which cells establish connections with their external environment. However, how changes in conformation and function of individual adhesion proteins regulate the dynamics of whole adhesion complexes remains largely elusive. This is because of the large separation in time and length scales between the dynamics of individual adhesion proteins (nanoseconds and nanometers) and the emergent dynamics of the whole adhesion complex (seconds and micrometers), and the limitations of molecular simulation approaches in extracting accurate free energies, conformational transitions, reaction mechanisms, and kinetic rates, that can inform mechanisms at the larger scales. In this review, we discuss models of integrin-based adhesion complexes and highlight their main findings regarding: (i) the conformational transitions of integrins at the molecular and macromolecular scales and (ii) the molecular clutch mechanism at the mesoscale. Lastly, we present unanswered questions in the field of modeling adhesions and propose new ideas for future exciting modeling opportunities.

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Introduction

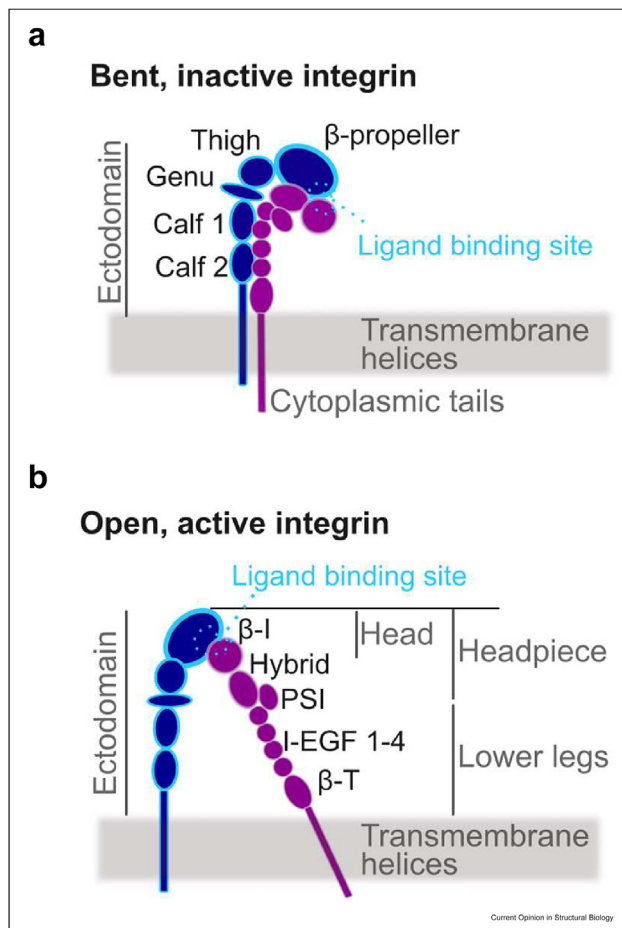
Adhesions between cells and the extracellular matrix (ECM) are important for physiological processes, including cell survival, growth, and proliferation, and for the maintenance of tissue-level structural integrity [1].

The integrin family of transmembrane receptors creates physical connections between cells and the ECM to transmit mechanical and allosteric signals that trigger downstream biochemical events and regulate cell function. The discovery of integrins occurred in the late 1970s and was followed by the identification of cytoplasmic integrin-associated proteins, including talin, vinculin, kindlin, α -actinin, and a multitude of signaling molecules. Together, integrins and associated proteins create highly dynamic adhesion complexes that vary composition and size over time, control ECM remodeling, and adapt their ECM binding strength to regulate cell function. Today, growing interest is directed toward understanding how exactly the individual proteins of the complex, including the integrins themselves, contribute to the constant remodeling of adhesion complexes.

Integrins are dimers composed of an α and β chain forming a large extracellular ectodomain, followed by two single-pass transmembrane helices and two short cytoplasmic tails connecting to cytoplasmic proteins and adaptors (Figure 1). The ectodomain of the α chain contains a large β -propeller domain, followed by the thigh domain, and two calf domains (Figure 1a). The ectodomain of the β chain consists of an N-terminal β -I domain, followed by the hybrid domain, the plexin–semaphorin–integrin domain (PSI), four cysteine-rich epidermal growth factor (EGF) modules (I-EGF 1–4) and β -T domain (Figure 1b). In the inactive, bent conformation, the headpiece is bent against the lower legs, and the site for binding an ECM ligand, which is at the interface between β -propeller and β -I domains, is oriented toward the cell membrane (Figure 1a). When integrin is active, its ectodomain is extended, with the ligand binding site oriented away from, rather than toward, the cell membrane (Figure 1b). The transition of integrin from bent to extended is associated with an increase in ligand binding affinity and the recruitment of intracellular adaptors, which leads to sequestration of more integrins and formation of adhesion complexes [2–4].

Molecular modeling work has focused on how atomistic interactions of integrin drive long-range motions during its transition from bent to extended conformations [5,6]. According to the switchblade model [7], the key pivots for integrin extension are the α chain genu, between the thigh and calf-1 domains, and the β chain knee, between I-EGF domains 1 and 2 (Figure 2).

Figure 1

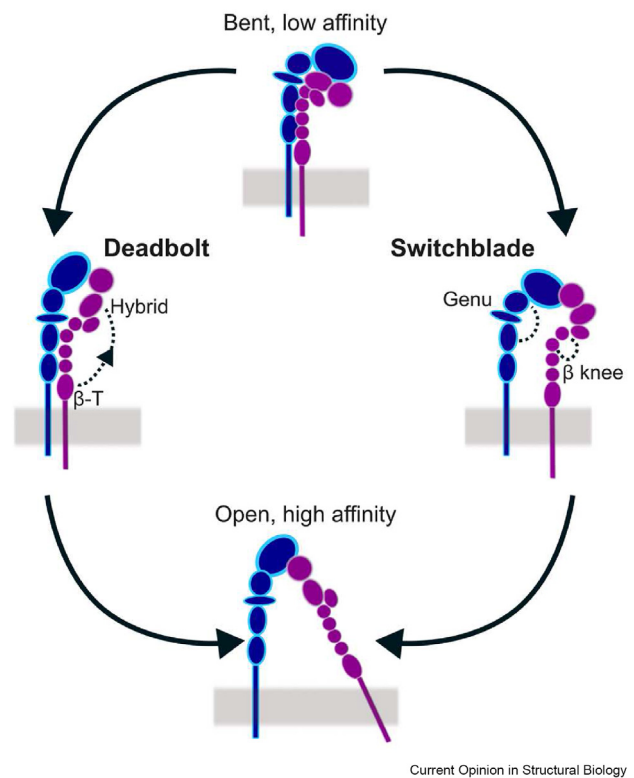


Schematic of integrin conformation and domain organization. Integrins transition between: **a**, a bent conformation that is inactive, with low affinity for binding external ligands; **b**, an open conformation that is active, with high affinity for binding external ligands. The different domains of the α chain (blue) and the β chain (purple) are indicated.

According to the deadbolt model [8], a hairpin loop in the β -T domain acts as a deadbolt that releases the β -I domain to extend integrin (Figure 2). These two models suggested two distinct pathways for integrin activation, but none of the current modeling approaches have fully captured these pathways because of the limited sampling of molecular simulations and the difficulty in choosing accurate collective variables for enhanced sampling methods. Mesoscale models of adhesion assembly have included experimental kinetic rates for integrin activation and deactivation and represented integrins as discrete particles switching between inactive and active conformations, without incorporating molecular details [9].

In cells, myosin-powered contractility and actin polymerization pushing against the cell edge drive a flow of

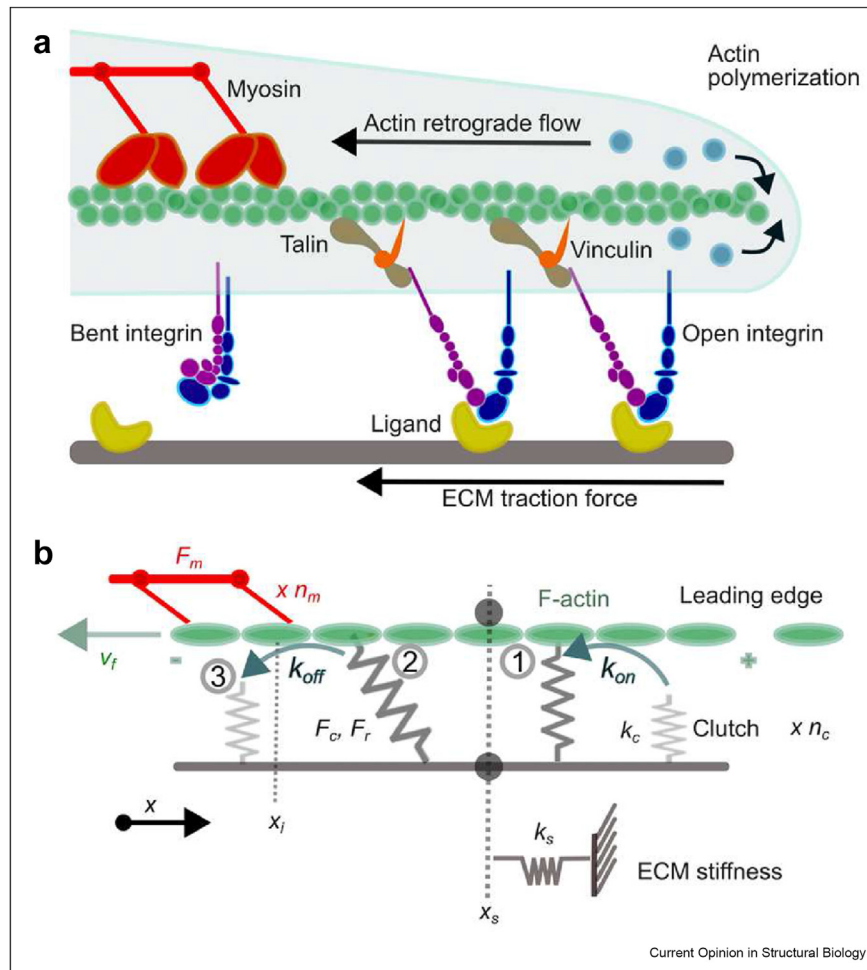
Figure 2



Models of integrin conformational activation. The transition from closed conformation to open conformation is characterized by intermediate integrin conformations, which differ in the degree of legs separation. According to the “deadbolt” (left) model, the headpiece extends before the legs separate laterally. According to the “switchblade” (right) model, the legs separate laterally before headpiece extension.

actin filaments that is transmitted as force to the ECM through adhesion adaptors and integrins (Figure 3). Similarly, ECM stiffness is transmitted through adhesions to the cytoskeleton. Internal and external forces change the composition of adhesions and their composition, size, and movement. A mesoscale computational model based on a coarse-grained description of adhesion complexes has described how cytoskeletal force controls the size of adhesions and physically remodel the ECM [10]. Analytical models incorporating multiple adhesion proteins into a minimal set of parameters have elucidated the relation between ECM stiffness and adhesion size, and how cytoskeletal force and adhesion velocity are connected, demonstrating a regime in which adhesion velocity depends linearly on the activation rate of integrins and ECM stiffness [11–13]. Continuum models integrating cytoskeleton architecture, adhesion adaptors, integrins and ECM stiffness have evaluated their relative contributions to the growth of adhesions [14]. In the early 2000s, the numerical implementation

Figure 3



An overview of the adhesion clutch model. a. New actin monomers (blue) are incorporated on to the barbed end of a pre-existing actin filament (green) facing the cell leading edge membrane. Myosin motors (red) exert contractile forces on the actin filament, pulling it away from the leading edge. Transmembrane integrin dimers (blue and purple) are either bent and freely diffusive or extended and bound to ECM ligands (yellow). If integrins are not engaged to connect actin to the ECM, then actin polymerization results in rapid retrograde actin flow, no net leading edge protrusion, and no traction force on the ECM. If integrins are engaged, then talin (dark yellow) and vinculin (orange) are also recruited and the forces generated by polymerization of the actin filament and myosin motors are physically transmitted to the ECM, resulting in slowing of actin retrograde flow, traction force on the ECM, and a net edge protrusion. The complex formed by ligand-bound integrins, talins and vinculins form the adhesion clutch. **b. Schematic representation of the motor-clutch model.** n_m myosin motors exert a maximum (a.k.a. stall) force, equal to $F_m n_m$, and pull an F-actin bundle, resulting in F-actin sliding at velocity v_f . This sliding is transmitted to the substrate through adhesions, which consist of n_c clutches. Each i th clutch can be engaged with F-actin (binding rate k_{on}) or disengaged (unbinding rate k_{off} that increases with force scaled to a rupture force F_r). Once engaged (1), the i th clutch is subjected to load from F-actin sliding (2) and eventually fails (3). F_c is the total force exerted by the engaged clutches, equal to $F_c = k_s x_s$, where k_s is the clutch stiffness and x_s is the final position of the substrate.

of a stochastic adhesion clutch model was carried out and its subsequent applications revealed the mechanisms by which adhesions sense ECM stiffness to mediate spreading and migration of cells [15–20]. The molecular clutch model is the current framework of reference for understanding how integrin-based adhesions govern cell function in response to internal and external stimuli.

In this review, we discuss some of the most important findings of computational approaches used for the study

of integrins and integrin-based adhesion complexes. We focus on methods that create the foundation to bridge the molecular dynamics of integrins to the scale of adhesion complexes. Lastly, we discuss unanswered questions in the field and present future modeling opportunities for bridging the gap between individual proteins and whole adhesion complexes.

All-atom and coarse-grained models of integrins

Many factors that regulate integrin conformational activation, such as ligand binding [21], divalent cations

[22], and mechanical force [23], are known from experiments. However, the short lifetimes of integrin conformational transitions (<1 s) relative to the temporal resolutions of live cell experiments (minutes) have made it challenging to capture the conformational pathway of integrin experimentally. To overcome these challenges, equilibrium and accelerated all-atom molecular dynamics (MD) simulations, combined with coarse-graining methods, are currently used.

Equilibrium and accelerated molecular dynamics simulations of integrin

Computational models based on cryo-EM or x-ray crystallographic structures of integrin have applied MD simulations to understand integrin dynamics with atomic details. These approaches rely on the application of empirical atomic force fields and use the classical Newtonian equations to calculate atomistic motions. Equilibrium MD simulations sample conformational ensembles around a local free energy minimum, and accelerated MD simulations sample a larger conformational space by decreasing the energy barriers between multiple minima. Types of accelerated MD simulations are steered MD, which applies force to perturb the atomic structure of a protein, and targeted MD, which biases an initial protein structure towards a known target, by imposing geometrical constraints.

Equilibrium, steered, and targeted MD simulations have typically used integrin $\alpha_v\beta_3$ or $\alpha_{IIb}\beta_3$. Steered MD simulations of integrins bound to talin or kindlin, and free energy perturbations of integrins bound to an external ligand, identified the molecular mechanisms underlining conformational activation, including affinity modifications and variations in the lifetimes of the ligand bound states [24–27].

MD-based studies identified the major energy barriers along the unbending pathway of integrin at the hybrid domain interfaces [28–30]. In the presence of an RGD sequence, which mimics ECM ligand binding, the hybrid domain opens because of molecular distortions involving reorientation of the α_7 and α_1 helices [28–30]. These distortions change the interaction surfaces of the hybrid domain with β propeller, β -I, PSI, β -T and EGF4 domains [28–31]. These changes, in turn, promote the swing out of the hybrid domain, followed by the opening of the hinge between the β -I and hybrid domains [32]. MD simulations also revealed that the opening of the β /hybrid hinge is inhibited when Ca^{2+} is bound instead of Mn^{2+} [29,33]. When force is applied to the ligand binding site, the opening of the β -I/hybrid hinge is accelerated [29,30], suggesting that the pathway for conformational activation is the same without and with force. Once in the extended conformation, an electrostatic interaction in the α chain genu stabilizes the conformation [31]. MD simulations also

demonstrated that, although these mechanisms are shared by $\alpha_v\beta_3$ and $\alpha_{IIb}\beta_3$ integrins, they can vary across integrin families [34].

Taken together, MD simulations provided important insights into the mechanisms underlying integrin conformational dynamics at the atomistic level. However, all-atom approaches are limited by high computational cost, low accuracy of empirical force fields, and limited sampling performance. MD simulations of integrin sample only up to tens of microseconds, a scale ~ 5 orders of magnitude shorter than the lifetime of conformational activation, and multiple independent replicates are often required to account for stochastic variability in initial conditions and for increasing the accuracy of the estimates of properties of interest, such as free energies, distribution of states, or kinetic parameters [35]. Additionally, empirical force fields inherently present low accuracy because they require predefined connectivity between atoms and are best suited to model a single protein conformation.

Coarse-grained simulations of integrin conformational activation

Capturing changes in integrin conformation with all-atom MD simulations remains a challenge because of the high computational cost, the need of multiple independent trajectories for improving sampling performance, and the low accuracy of empirical force fields. These challenges limit the access of all-atom MD simulations to collective global motions at the macromolecular scale. To address these challenges, bottom-up coarse-graining (CG) methods have provided realistic approximations of integrin macromolecular dynamics by grouping atoms into representative beads and incorporating bond breakage and formation to decrease the barriers across multiple conformations. Results identified the main directions of integrin global motions at the onset of conformational activation and the effects of residue mutations on these motions.

Bottom-up elastic and mixed elastic/plastic network models (ENMs) represent integrins as networks of beads interconnected by harmonic or mixed harmonic and anharmonic potentials, respectively [36–39]. Analysis of the normal modes of an ENM model of the bent $\alpha_v\beta_3$ integrin ectodomain has demonstrated that weakening the interactions between hybrid and β -T domains promotes the displacement of the hybrid domain away from the β leg, such that further long-range conformational changes become easier [36]. Similarly, weakening the interactions between ENM beads based on MD analysis has promoted integrin motions along a frustrated energy landscape, characterized by leg separation, headpiece extension and flattening of the α and β chain knees [38,39]. MD-based CG models of several activating mutants have revealed the existence of

mechanosensitive mutants, able to lower the force required for conformational activation [37]. The result that molecular changes in integrin confer different responses to force has been further confirmed by MD simulations on $\alpha_{IIb}\beta_3$ [34].

Taken together, bottom-up CG models of integrin have provided a deeper understanding of the long-range motions occurring during integrin conformational activation. However, the full pathway of integrin conformational activation remains outside the reach of current methods.

Adhesion clutch models

Mesoscale models of adhesion complexes have simplified adhesions as mechanical clutches between cells and the ECM. Adhesion clutch models have focused on the mechanisms by which adhesions convert actomyosin contractility and actin retrograde flow into ECM traction [40]. In a typical adhesion-clutch simulation, clutches have elastic properties and initially bind the ECM with a defined first-order rate constant to transmit the motion of the cytoskeleton to the ECM. As cytoskeletal force builds on the ECM-bound clutches (Figure 3), actin flow slows and ECM traction force increases. The clutches eventually fail and release the tension, which corresponds to an increase in actin flow speed and a decrease in ECM traction force. Then, the cycle starts again, leading to a “load-and-fail” behavior [15,17].

The adhesion clutch mechanism has been implemented in stochastic particle-based simulations, finite elements approaches, and analytical theory [17,41–43]. The adhesion clutch model has found broad application to simulate cell behaviors due to clutch heterogeneity [19], mechanical heterogeneity and anisotropy in 1D and 2D nanofibrous environments [44], varying adhesion levels in brain tissue [45], perturbation by anti-cancer drugs [46], stress-relaxing environments [42], stress-relaxing cytoskeletal dynamics [47], RhoGTP signaling [48], 2D and 3D cell migration [49,50], and negative durotaxis, i.e., migration toward softer environments [16].

Adhesion clutch models have also incorporated the effects of talin and vinculin as adaptors between integrins and actin filaments (Figure 3), using experimentally extracted kinetic rate constants [51]. Further modeling work based on Brownian dynamics implemented the catch bond kinetics of individual integrin-ECM bonds, again using experimental lifetime versus force relations of the bonds [52]. These studies, collectively, have demonstrated that adhesion clutches aid in membrane protrusion depending on the proteins involved and bond dynamics, regulate force transmission, mechanosensing and filopodia dynamics, and can also govern cycles of membrane oscillations.

Given its ability to explain a large variety of cell phenomena, the adhesion clutch model is the framework of reference to understand adhesion-mediated processes in cells. However, given the complexity of the dynamic changes in conformation and function of individual adhesion proteins in the clutch, the exact relation between the dynamics of individual protein components and the clutch mechanism remains largely unexplored. It remains a major challenge to convert atom-level structural information into parameter values, such as clutch stiffness, on- and off-rate constants, and characteristic catch and bond rupture forces, that inform the motor-clutch model at the mesoscale.

Concluding remarks

Cells in almost any physiological setting, from neurons within the brain to immune cells crossing tissues, form adhesions with the ECM. The development and application of MD simulations and CG models for the study of adhesion proteins, combined with mesoscale and continuum models of whole adhesion complexes, have allowed the field to gain a fundamental understanding of how cell-ECM adhesions operate under internal and external stimuli. MD and CG methods have revealed the molecular underpinnings for integrins conformational activation, which are inaccessible by current experimental approaches. Mesoscale and continuum models of adhesion clutches have revealed how integrin-based adhesions regulate force transmission, cell edge protrusion, cell mechanosensing, and migration by mediating internal and external force transmission. However, these models are still used in isolation. Because of the lack of experimental approaches with sufficient spatial and temporal resolution to monitor how the changes in structure and function of multiple adhesion proteins govern the dynamics of whole adhesion complexes, it will be interesting in the future to develop multiscale models that rigorously relate molecular mechanisms with mesoscale properties of adhesions and the clutch mechanism.

Today, while this effort is ongoing, several outstanding questions remain to be addressed. At the molecular scale, how force is transmitted and distributed across adhesion proteins remains largely elusive. Addressing this question will require developing computational methods able to track force distribution across protein domains and residues, while simulating multiple interacting proteins. Second, how exactly the emergent global motions of adhesion proteins emerge from atomistic rearrangements remains largely elusive. Developing new methods able to directly access molecular events such as rearrangements of secondary structure elements, changes in interaction surfaces between domains, and their relationship with long-range structural motion will allow us to relate molecular with macromolecular dynamics of integrins and other adhesion proteins. For these studies, multiple atomistic MD

replica will be required to decrease the error of the average of the estimates and infer functional mechanisms. Lastly, it is highly likely that atomistic changes and long-range global motions of integrins and other adhesion proteins are correlated with the assembly and stabilization of whole adhesion complexes, and with the efficiency of force transmission to and from the ECM. Conformational and affinity changes of adhesion proteins can also affect the movement, size, and lifetimes of the whole adhesion. Developing new multiscale methods exploring relationships between the molecular and macromolecular properties of adhesion proteins and the mesoscale properties of adhesion complexes will allow the field to address these open questions. Multiscale methods for the study of adhesion complexes will further enable the testing and refining of novel clutch models to elucidate how the clutch mechanisms more precisely mediate cell function. Addressing these and other open questions is thus likely to lead to exciting new model developments in the coming years.

Conflict of interest statement

Nothing declared.

Data availability

No data were used for the research described in the article.

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