Evaluation of Time Integration Schemes for the Generalized Interpolation Material Point Method

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Recurring Themes in Time Integration

- Close coupling between space and time
- Low order is better
- Linear error theories not applicable
- Behavior drastically different for large deformation
- So what's really going on?
- Manufactured solutions measure effect of time integration on
 - Accuracy
 - Stability

Compare USF and USL

Update Stress First

Update Stress Last

$$\nabla \mathbf{v}_p^n = \sum_i \mathbf{v}_i \mathbf{G}_{ip}$$

$$\mathbf{F}_{p}^{n+1} = \mathbf{F}_{p}^{n} + \nabla \mathbf{v}_{p}^{n} \cdot \mathbf{F}_{p}^{n} \Delta t$$
$$\mathbf{a}_{i} \left(\mathbf{\sigma} \left(\mathbf{F}_{p}^{n+1} \right) \right)$$

$$\mathbf{a}_i \left(\mathbf{\sigma} \left(\mathbf{F}_p^n \right) \right)$$

$$\nabla \mathbf{v}_p^{n+1} = \sum_i (\mathbf{v}_i + \mathbf{a}_i \Delta t) \mathbf{G}_{ip}$$

 $\mathbf{F}_{p}^{n+1} = \mathbf{F}_{p}^{n} + \nabla \mathbf{v}_{p}^{n+1} \cdot \mathbf{F}_{p}^{n} \Delta t$

Compare Centered-Difference and USL

$$\mathbf{v}_i^{n-1/2} = \frac{\sum \mathbf{v}_p^{n-1/2} m_p \mathbf{S}_{ip}}{\sum m_p \mathbf{S}_{ip}}$$

$$\mathbf{v}_i = \frac{\sum \mathbf{v}_p^n m_p \mathbf{S}_{ip}}{\sum m_p \mathbf{S}_{ip}}$$

$$\mathbf{v}_{p}^{n+1/2} = \mathbf{v}_{p}^{n-1/2} + \mathbf{a}_{p}^{n} \Delta t$$
$$\mathbf{x}_{p}^{n+1} = \mathbf{x}_{p}^{n} + \mathbf{v}_{p}^{n+1/2} \Delta t$$
$$\mathbf{F}_{p}^{n+1} = \mathbf{F}_{p}^{n} + \nabla \mathbf{v}_{p}^{n+1/2} \cdot \mathbf{F}_{p}^{n} \Delta t$$

$$\mathbf{v}_{p}^{n+1} = \mathbf{v}_{p}^{n} + \mathbf{a}_{p} \Delta t$$
$$\mathbf{x}_{p}^{n+1} = \mathbf{x}_{p}^{n} + \mathbf{v}_{p} \Delta t$$
$$\mathbf{F}_{p}^{n+1} = \mathbf{F}_{p}^{n} + \nabla \mathbf{v}_{p}^{n+1} \cdot \mathbf{F}_{p}^{n} \Delta$$

Initialization to a Negative Half Time Step

If you know the answer:
$$v_p = v(t = -k/2)$$

Use data at time = 0:
$$v_p^{-1/2} = v_p^0 - \frac{\Delta t}{2} a_p^0$$

The easy way: if first time step then

$$a_i^0 = \frac{1}{2}a_i^0$$

Axis-Aligned Displacement in a Unit Square

$$u = \begin{pmatrix} A\sin(\pi X)\cos(c\pi t) \\ A\sin(\pi Y)\sin(c\pi t) \\ 0 \end{pmatrix}$$



Functions of coordinate directions only

Corners and edges of GIMP particles remain aligned

Sliding boundaries – zero normal velocity at surface.

Axis-Aligned Displacement – cont.

Diagonal terms only:

$$F = \begin{bmatrix} 1 + A\pi \cos(\pi X) \cos(c\pi t) & 0 & 0 \\ 0 & 1 + A\pi \cos(\pi Y) \sin(c\pi t) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Stress

Momentum

 $\mathbf{P} = \lambda \ln(J)\mathbf{F}^{-1} + \mu \mathbf{F}^{-1} (\mathbf{F}\mathbf{F}^{T} - \mathbf{I}) \qquad \nabla \cdot \mathbf{P} + \rho_{0}\mathbf{b} = \rho_{0}\mathbf{a}$

Solve for body force from momentum:

$$\mathbf{b} = \frac{\pi^2}{\rho_0} \begin{pmatrix} u_X \left[\lambda \left(1 - \ln(F_{XX} F_{YY}) \right) F_{XX}^{-2} + \mu \left(1 + F_{XX}^{-2} \right) - E \right] \\ u_Y \left[\lambda \left(1 - \ln(F_{XX} F_{YY}) \right) F_{YY}^{-2} + \mu \left(1 + F_{YY}^{-2} \right) - E \right] \\ 0 \end{pmatrix}$$

Axis-Aligned Displacement



Definition of Error at a Particle

$$\delta = \left(x_p^n - X_p^n\right) - u\left(X_p^n, t^n\right)_{EXACT}$$

For a smooth problem in space and time check all particles and all time steps:

$$L_{\infty} = \max(\delta)$$

Spatial Convergence



CD-GIMP is 2nd order – the initialization shortcut works USL changes to 1st order – effect of half step initialization

Minor change to UGIMP causes large error

USF and MPM visually OK but poor accuracy



Most methods display zero temporal convergence until stability is lost, even though CD-GIMP is formally 2nd order in time.

USL loses one spatial order, and gains one temporal – sum of 2?

We conclude that spatial error dominates temporal error such that reduced CFL has no benefit.

GIMP Convergence Order

Hypothesis

GIMP is second order in space when:

- 1. Problem is smooth in space and time
- 2. Particle edges are aligned
- 3. Material boundary is accurately represented

Therefore the 2D code is verified.

Expanding Ring

Free surfaces with implied zero normal stress GIMP particle edges not aligned More general and representative



Expanding Ring – Displacement

Radial Symmetry
$$u(R,t) = T(t)[c_3R^3 + c_2R^2 + c_1R]$$

Capital "R" is radius in reference configuration.

X and Y Displacement Components:

$$\mathbf{u} = \begin{bmatrix} T[c_3R^3 + c_2R^2 + c_1R]\cos(\theta) \\ T[c_3R^3 + c_2R^2 + c_1R]\sin(\theta) \\ 0 \end{bmatrix}$$



radius

Expanding Ring: Cartesian Coordinates

Gradient Operators in terms of R and θ :

 $\mathbf{F} = \mathbf{I} + \begin{bmatrix} \frac{\partial u_X}{\partial X} & \frac{\partial u_X}{\partial Y} \\ \frac{\partial u_Y}{\partial X} & \frac{\partial u_Y}{\partial Y} \\ 0 & 0 \end{bmatrix}$

$$\frac{\partial f(R,\theta)}{\partial X} = \frac{\partial f}{\partial R} \cos(\theta) - \frac{\partial f}{\partial \theta} \frac{\sin(\theta)}{R}$$
$$\frac{\partial f(R,\theta)}{\partial Y} = \frac{\partial f}{\partial R} \sin(\theta) + \frac{\partial f}{\partial \theta} \frac{\cos(\theta)}{R}$$
$$0$$
Stress with zero Poisson's ratio
$$\mathbf{P} = \mu \mathbf{F}^{-1} \left(\mathbf{F} \mathbf{F}^{T} - \mathbf{I} \right)$$

Now we let Maple do the hard part . . .

Stress Matrix

 $\frac{1}{2}ET(T^{2}\cos(H)^{2}c^{2}Rc^{2}l^{2}+3T^{2}c^{2}R^{2}\cos(H)^{2}c^{2}l+2T^{2}\cos(H)^{2}c^{3}R^{2}c^{2}l^{2}+2T\cos(H)^{2}c^{2}Rc^{2}l+4T\cos(H)^{2}c^{3}R^{2}c^{2}l^{2}+2T\cos(H)^{2}c^{2}Rc^{2}l^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}c^{2}Rc^{2}l^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2T\cos(H)^{2}+2Tc^{2}+2T\cos(H)^{2}+2Tc^{2}+2T\cos(H)^{2}+2Tc^{2}+2T\cos(H)^{2}+2Tc^{2}+2T\cos(H)^{2}+2Tc^{2}+2T\cos(H)^{2}+2Tc^$ $+ 10 T^{2} c^{3} R^{3} \cos(H)^{2} c^{2} c^{1} + 8 T^{2} c^{3} R^{4} \cos(H)^{2} c^{1} + 13 T^{2} c^{3} R^{5} \cos(H)^{2} c^{2} + 9 T^{2} c^{3} R^{4} \cos(H)^{2} c^{2} + 10 T \cos(H)^{2} c^{3} R^{3} c^{2}$ $+5 T c 2^{2} R^{2} + 7 T c 3^{2} R^{4} + T^{2} c 1^{3} + 3 T c 1^{2} + 2 T^{2} c 2^{3} R^{3} + 3 T^{2} c 3^{3} R^{6} + 8 T c 2 R c 1 + 10 T c 3 R^{2} c 1 + 12 T c 3 R^{3} c 2 + 7 T^{2} c 3 R^{4} c 2^{2}$ $+2T^{2}c^{3}R^{3}\cos(H)^{2}+2c^{2}R+3T\cos(H)^{2}c^{2}R^{2}+8T\cos(H)^{2}c^{3}R^{4}+2\cos(H)^{2}c^{2}R+4\cos(H)^{2}c^{3}R^{2})/(6R^{2}+1)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{4}+2\cos(H)^{2}c^{2}R+4\cos(H)^{2}c^{3}R^{2})/(6R^{2}+1)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{4}+2\cos(H)^{2}c^{3}R^{4}+2\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{4}+2\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}c^{3}R^{2}+8T\cos(H)^{2}+$ $1 + 3T^{2}c^{3}R^{4} + 5T^{2}c^{3}R^{3}c^{2} + 4T^{2}c^{3}R^{2}c^{2} + 3T^{2}c^{2}Rc^{2} + 3T^{2}c^{2}R^{2} + 3Tc^{2}R^{2} + 3Tc^{2}R^{2} + 3Tc^{2}R^{2}), \frac{1}{2}ET\cos(H)\sin(H)R$ $(2 c 3 R + c 2) (2 + 3 T^{2} c 3^{2} R^{4} + 5 T^{2} c 3 R^{3} c 2 + 4 T^{2} c 3 R^{2} c l + 3 T^{2} c 2 R c l + T^{2} c l^{2} + 2 T c l + 2 T^{2} c 2^{2} R^{2} + 3 T c 2 R + 4 T c 3 R^{2}) / (2 c 3 R + c 2) (2 + 3 T^{2} c 3^{2} R^{4} + 5 T^{2} c 3 R^{3} c 2 + 4 T^{2} c 3 R^{2} c l + 3 T^{2} c 2 R c l + T^{2} c l^{2} + 2 T c l + 2 T^{2} c 2^{2} R^{2} + 3 T c 2 R + 4 T c 3 R^{2}) / (2 c 3 R + c 2) (2 + 3 T^{2} c 3^{2} R^{4} + 5 T^{2} c 3 R^{3} c 2 + 4 T^{2} c 3 R^{2} c l + 3 T^{2} c 2 R c l + T^{2} c l^{2} + 2 T c l + 2 T^{2} c 2^{2} R^{2} + 3 T c 2 R + 4 T c 3 R^{2}) / (2 c 3 R + c 2) (2 + 3 T^{2} c 3^{2} R^{4} + 5 T^{2} c 3 R^{3} c 2 + 4 T^{2} c 3 R^{2} c l + 3 T^{2} c 2 R c l + T^{2} c l^{2} + 2 T c l + 2 T^{2} c 2^{2} R^{2} + 3 T c 2 R + 4 T c 3 R^{2}) / (2 c 3 R c l + 1 C c 3 R^{2} c l + 3 T^{2} c 2 R c l + 1 C c 2 R^{2} c l + 3 T^{2} c l + 3$ $1 + 3 T^{2} c^{3} R^{4} + 5 T^{2} c^{3} R^{3} c^{2} + 4 T^{2} c^{3} R^{2} c^{1} + 3 T^{2} c^{2} R c^{1} + T^{2} c^{2} r^{2} + 2 T c^{1} + 2 T^{2} c^{2} R^{2} + 3 T c^{2} R + 4 T c^{3} R^{2}), 0$ $\frac{1}{2}ET\cos(H)\sin(H)R(2c3R+c2)$ $(2+3T^{2}c^{3}R^{4}+5T^{2}c^{3}R^{3}c^{2}+4T^{2}c^{3}R^{2}c^{1}+3T^{2}c^{2}Rc^{1}+T^{2}c^{2}r^{2}+2Tc^{1}+2T^{2}c^{2}R^{2}+3Tc^{2}R+4Tc^{3}R^{2})/((2+3T^{2}c^{2}R^{2}+3T^{2}+3T^{2}c^{2}R^{2}+3T^{2}+3T^{2}c^{2}+3T^{2}+$ $1 + 3T^{2}c^{3}R^{4} + 5T^{2}c^{3}R^{3}c^{2} + 4T^{2}c^{3}R^{2}c^{2} + 3T^{2}c^{2}Rc^{2} + 2Tc^{2}c^{2}R^{2} + 3Tc^{2}Rc^{2} + 4Tc^{3}R^{2}), -\frac{1}{2}TE(1 + 2T^{2}c^{2}R^{2} + 3Tc^{2}Rc^{2} + 4Tc^{3}R^{2}), -\frac{1}{2}TE(1 + 2T$ $T^{2}\cos(H)^{2}c^{2}Rc^{2}l^{2} + 3T^{2}c^{2}R^{2}\cos(H)^{2}c^{2}l + 2T^{2}\cos(H)^{2}c^{3}R^{2}c^{2}l^{2} + 2T\cos(H)^{2}c^{2}Rc^{2}l + 4T\cos(H)^{2}c^{3}R^{2}c^{2}l^{2}$ $+ 10 T^{2} c^{3} R^{3} \cos(H)^{2} c^{2} c^{1} + 8 T^{2} c^{3}^{2} R^{4} \cos(H)^{2} c^{1} + 13 T^{2} c^{3}^{2} R^{5} \cos(H)^{2} c^{2} + 9 T^{2} c^{3} R^{4} \cos(H)^{2} c^{2} + 10 T \cos(H)^{2} c^{3} R^{3} c^{2}$ $-8 T c 2^{2} R^{2} - 15 T c 3^{2} R^{4} - T^{2} c 1^{3} - 3 T c 1^{2} - 4 T^{2} c 2^{3} R^{3} - 9 T^{2} c 3^{3} R^{6} - 10 T c 2 R c 1 - 14 T c 3 R^{2} c 1 - 22 T c 3 R^{3} c 2 - 16 T^{2} c 3 R^{4} c 2^{2}$ $-7 T^{2} c^{3} R^{2} cl^{2} - 15 T^{2} c^{3}^{2} R^{4} cl - 21 T^{2} c^{3}^{2} R^{5} c^{2} - 8 T^{2} c^{2}^{2} R^{2} cl - 5 T^{2} c^{2} R cl^{2} - 22 T^{2} c^{3} R^{3} c^{2} cl + 6 T^{2} c^{3}^{3} R^{6} \cos(H)^{2} - 6 c^{3} R^{2} cl^{2} R^{2} cl^{2} - 22 T^{2} c^{3} R^{3} c^{2} cl + 6 T^{2} c^{3}^{3} R^{6} \cos(H)^{2} - 6 c^{3} R^{2} cl^{2} - 22 T^{2} c^{3} R^{3} c^{2} cl^{2} + 6 T^{2} c^{3}^{3} R^{6} cos(H)^{2} - 6 c^{3} R^{2} cl^{2} - 22 T^{2} c^{3} R^{3} c^{2} cl^{2} + 6 T^{2} c^{3} r^{3} c^{2} cl^{2} - 22 T^{2} c^{3} r^{3} cl^{2} - 22 T^{2} c^{3} r^{3} cl^{2} - 22 T^{2} c^{3} r^{3} cl^{2} cl^{2} - 22 T^{2} cl^{2} cl^{2} - 22 T^{2} cl^{2} cl^{2} cl^{2} - 22 T^{2} cl^{2} cl^{2} cl^{2} - 22 T^{2} cl^{2} cl^{2} cl^{2} cl^{2} - 22 T^{2} cl^{2} cl^{2} cl^{2} cl^{2} - 22 T^{2} cl^{2} cl^{2}$ $-2c1 + 2T^{2}c2^{3}R^{3}\cos(H)^{2} - 4c2R + 3T\cos(H)^{2}c2^{2}R^{2} + 8T\cos(H)^{2}c3^{2}R^{4} + 2\cos(H)^{2}c2R + 4\cos(H)^{2}c3R^{2}) / (1 + 2T^{2}c2^{3}R^{3}\cos(H)^{2} - 4c2R + 3T\cos(H)^{2}c2^{2}R^{2} + 8T\cos(H)^{2}c3^{2}R^{4} + 2\cos(H)^{2}c2R + 4\cos(H)^{2}c3R^{2}) / (1 + 2T^{2}c2^{3}R^{3}) / (1$ $1 + 3T^{2}c^{3}R^{4} + 5T^{2}c^{3}R^{3}c^{2} + 4T^{2}c^{3}R^{2}c^{1} + 3T^{2}c^{2}Rc^{1} + T^{2}c^{2}r^{2} + 2Tc^{1} + 2T^{2}c^{2}R^{2} + 3Tc^{2}R + 4Tc^{3}R^{2}), 0$

[0, 0, 0]

P :=

Find c_1 , c_2 , and c_3 by rotating the stress matrix

Rotation Matrix Q:
$$\mathbf{Q} = \begin{bmatrix} \cos(\theta) & \sin(\theta) & 0 \\ -\sin(\theta) & \cos(\theta) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$
$$\mathbf{P'} = \mathbf{Q} \mathbf{P} \mathbf{Q}^{T} = \begin{bmatrix} \alpha & 0 & 0 \\ 0 & \beta & 0 \\ 0 & 0 & 0 \end{bmatrix} \qquad \begin{aligned} \alpha(R_{o}) = 0 \\ \alpha(R_{I}) = 0 \\ u(R_{o}) = T \end{aligned}$$

Maple finds a simple answer:

$$\{c3 = 2\frac{1}{Ro^{2}(3Ri - Ro)}, c2 = -3\frac{Ro + Ri}{Ro^{2}(3Ri - Ro)}, c1 = 6\frac{Ri}{(3Ri - Ro)Ro}\},\$$

 $u(R_0) = T$

Solve for Body Force

$$T(t) = A\cos\left(\sqrt{\frac{E}{\rho_0}} \pi t\right)$$

Maple generates C-compatible code for b

Expanding Ring Results



Normal Stress

Ring – Spatial Convergence



The miss-alignment of particles and the stair-stepped surface now dominate the error.

UGIMP gives up only for the highest resolutions.



UGIMP and GIMP nearly same accuracy

Same trends otherwise: little temporal convergence

Conclusions

The Method of Manufactured solutions allows order of accuracy to be demonstrated for realistic large deformation problems.

The CD-GIMP combination is significantly better than choices involving UGIMP, MPM, USF, and USL.

Formal temporal orders of accuracy are usually not observed in real solutions because spatial error dominates temporal.

CD-GIMP can be 2nd order if problem is smooth, particle edges are aligned, and surface is well-represented. Convergence drops to 1st order for non-aligned particles.

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Expanding Disk: C code from Maple